






# Load-Independent Push–Pull Class E<sup>2</sup> Topology With Coupled Inductors for MHz-WPT Applications

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**Abstract**—This article investigates the modeling and implementation of the load-independent operation for the push–pull Class E topology with coupled inductors. The analytical modeling is first presented and provides a new design methodology, which distantly improves the soft-switching performance, load regulation, and efficiency compared with the conventional inductor-coupled Class E topology. An LCC-S resonant Class E<sup>2</sup> converter, which combines the load-independent inductor-coupled inverter and active rectifier, is proposed for MHz-wireless power transfer (MHz-WPT) applications. The active inductor-coupled Class E rectifier maintains soft switching, constant voltage gain, and zero-phase-angle input at various loads. A three-winding magnetic structure is also proposed to improve the coupled inductors by enhanced self-inductance, which features more effective harmonic suppression with lower magnetic losses and volumes. By the proposed methodology and the improved magnetic structure, the converter features negligible harmonic contents and maintains soft switching with a constant voltage conversion ratio over the entire load range. A 6.78-MHz WPT prototype was built to verify the proposed methodology with detailed parameter design and hardware implementation. The prototype realizes load-independent zero-voltage switching and maintains an approximately constant voltage gain with +5/−4% variance from no-load to 320-W output power. The measured dc–dc efficiency reaches 89.3% at 210 W.

**Index Terms**—Active rectifier, Class E, coupled inductor, load-independent, wireless power transfer (WPT).

## I. INTRODUCTION

WIRELESS power transfer (WPT) based on the magnetic coupling is increasingly used in various applications,

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including portable devices, medical implants, unmanned aerial vehicles (UAV), and electric vehicles [1]–[5]. With the development of wide-bandgap (WBG) devices, MHz-WPT operating in industrial, scientific, and medical (ISM) band is becoming promising to improve system efficiency with volumetric and gravimetric power density, compared with most existing WPT systems operating at several hundred kHz.

In MHz-WPT systems, the inverters and rectifiers are to generate and rectify the high-frequency power transferred by the contactless magnetic resonant tanks. As the operating frequency increases to the multi-MHz range, soft-switching topologies and high-quality-factor magnetic components are crucial for systems' efficiency and reliability. The active full-bridge and half-bridge topologies have been widely used in WPT systems. They require fewer magnetic components and are conducive to realize higher power density [6], [7]. However, since the switching period is substantially shortened in the multi-MHz range, it requires complex gate driver control to perform reliable zero-voltage switching (ZVS) for the active bridges applied with variable load.

To avoid the inherent issues of active bridges, the resonant Class E topology and its derivative forms are preferred topologies for MHz-WPT systems [8]–[12]. The conventional Class E inverters comprise a low-side switch that can easily operate in the multi-MHz range. Generally, the switch realizes both zero-voltage and zero-voltage-derivative turn-ON at the optimal load. However, the soft-switching operation is sensitive to load variation. Although feedback control is a potential solution to adjust the switching behavior at various loads, it requires a complex design for reliable operation [11], [12].

The soft switching performance can be improved by using finite dc-feed inductance in open-loop Class E topologies with variable load. Roslaniec *et al.* [13] implement an experiment-based approach to realize nearly invariant ZVS operation over a wide load range. Aldhaher *et al.* [14] introduce the load-independent single-switch Class E topology into MHz applications. In theory, the load-independent operation achieves perfect ZVS and constant voltage gain regardless of load variation. Nevertheless, the much lower dc-feed inductance results in dramatically increase of harmonic currents on dc sides. The harmonics will cause not only substantial losses but also electromagnetic interference (EMI) issues.

Interleaving approaches can be implemented to noticeably reduce the harmonic contents [15]–[21]. By using two switching units operating in push–pull mode, the fundamental and

odd-order harmonics circulate within the units, while only the even-order harmonics flow into dc sides. The design process is basically the same as single-switch inverters/rectifiers. Liu and Chen [18] use an enhanced push-pull Class E inverter to power multiple reconfigurable receivers of dual-band WPT systems. In [19], the active push-pull Class E rectifier achieves the load-independent operation and has zero-phase-angle (ZPA) input over the load range. Gu *et al.* [20] propose a push-pull Class  $\Phi_2$  inverter, which features very low input current ripple and the load-independent ZVS. And, the push-pull  $\Phi_2$  was implemented in a 1.7-KW 6.78-MHz WPT system [21].

It is noticed that the discrete magnetic components dominate the size of the power stage, while also limiting the further power density improvement of the system. Air-core or low-permeability-core inductors usually perform high-quality-factors in the MHz region, but it is challenging to implement compact inductors for Class E inverters/rectifiers [7], [22]–[26].

An investigation has been reported in [27] to integrate and couple the inductors of the push-pull Class E inverters. The coupled inductor was also implemented for active rectification [28]. This integration reduces magnetic component count and volumes. Moreover, the coupled inductors perform noticeable suppression of even-order harmonics on dc sides. Nevertheless, the existing design cannot achieve the load-independent soft-switching, which is even critical for MHz-WPT systems. The involved magnetic coupling breaks the previous circuit model achieving load-independent operation.

This article investigates the modeling, implementation, and improvement of the inductor-coupled push-pull Class E topology achieving load-independent operation, which realizes soft switching, low-harmonic contents, and constant voltage gain over the entire load range. Section II presents the deduction and obtains a unique parameter set that achieves the load-independent operation, which distantly improves the soft-switching performance and load regulation compared with the conventional inductor-coupled Class E topology in [27]. Section III discusses the active coupled-inductor Class E rectifier achieving ideal ZVS and ZPA input with variable load. An *LCC-S* resonant Class E<sup>2</sup> topology is thus proposed to achieve load-independent soft-switching and constant output voltage over the entire load range. Furthermore, an improved magnetic structure is proposed to implement the coupled inductors with enhanced self-inductance, which performs much better harmonic suppression with lower losses and volumes than the previous work. In Section IV, we built a 6.78-MHz WPT prototype to verify the proposed methodology with detailed parameter design and hardware implement. The experimental result is discussed and compared with our previous work in Section V. It shows expected soft-switching performance, low-harmonic contents, and efficiency improvement. Finally, Section VI concludes this article.

## II. MODELING OF LOAD-INDEPENDENT OPERATION FOR INDUCTOR-COUPLED PUSH-PULL CLASS E INVERTERS

The load-independent single-switch Class E topology has been verified and implemented [14], [29]. Nevertheless, the involved magnetic coupling breaks the previous circuit model.

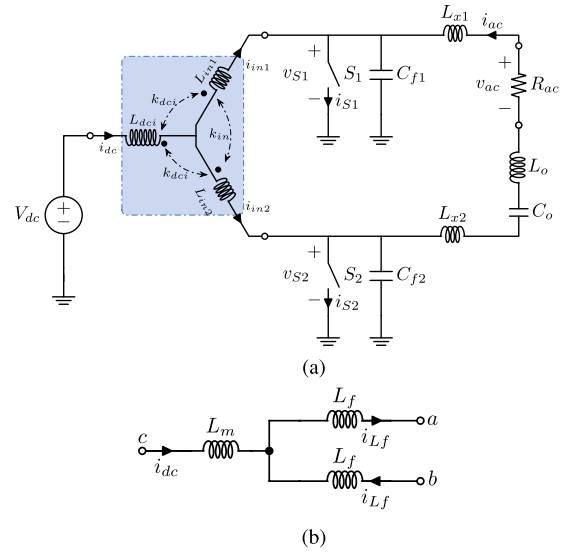


Fig. 1. (a) Push-pull Class E inverter with three-winding coupled inductor. (b) Equivalent decoupled circuit of coupled inductor. The output filter ( $L_o$  and  $C_o$ ) resonates at the operation frequency.  $L_x$  is the additional inductance, which can be implemented by individual inductors or merges with  $L_o$ .

Thus, in this section, we derive the new circuit operating model for the inductor-coupled push-pull Class E topology.

### A. Circuit Analysis

Fig. 1 illustrates the inductor-coupled push-pull Class E inverter. The proposed topology comprises a three-winding magnetic structure and additional inductance  $L_x$ , both of which have not been discussed in the previous work [27]. The additional inductance causes inherent differences in circuit operation since it changes the resonant network. To simplify the analysis, assumptions are made as follows.

1) The switches have zero ON-resistance and infinite OFF-resistance. The passive components are lossless.

2)  $S_1$  is ON during  $0 < \omega t \leq \pi$  and OFF during  $\pi < \omega t \leq 2\pi$ , where  $\omega$  is the angular velocity. In contrast,  $S_2$  is OFF during  $0 < \omega t \leq \pi$  and ON during  $\pi < \omega t \leq 2\pi$ . Since the switches operate at differential mode, the switch voltage can be defined as  $v_S(\omega t) = v_{S1}(\omega t) = v_{S2}(\omega t + \pi)$ .

3) The coupled inductor has large enough self-inductance, which makes the current ripples of  $i_{dc}$  negligible.

4) The characteristic impedance  $\sqrt{L_o/C_o}$  of the output filter is high enough to form a sinusoidal output current  $i_{ac}$  with phase  $\varphi$  and amplitude  $I_{ac}$ , i.e.,

$$i_{ac}(\omega t) = I_{ac} \cdot \sin(\omega t + \varphi). \quad (1)$$

As in Fig. 1(b), the three-winding coupled inductor can be decoupled to the mutual-inductance  $L_m$  and the leakage-inductance  $L_f$ , of which modeling is presented in Section III. This magnetic structure is proposed to enhance the equivalent self-inductance, which dominates the harmonic suppression. With proper winding arrangement, the inductance satisfies  $L_m \gg L_f$ , and thus the current ripples of  $i_{dc}$  are eliminated. Hence, the winding currents can be defined as  $i_{L_f}(\omega t) = i_{in1}(\omega t) = i_{in2}(\omega t + \pi)$ . Then,  $i_{L_f}$  during  $\pi < \omega t \leq 2\pi$  can be

expressed by

$$i_{L_f}(\omega t) = I_{L_f\pi} - \frac{1}{\omega L_e} \int_{\pi}^{\omega t} v_{S_n}(\omega t_1) d\omega t_1 \quad (2)$$

where  $L_e = 2L_f$ , and  $I_{L_f\pi}$  is the initial value of  $I_{L_f}$  at  $\omega t = \pi$ . Thus, the OFF-state switch voltage for  $\pi < \omega t \leq 2\pi$  satisfies

$$\omega C_f \cdot \frac{dv_S(\omega t)}{d\omega t} = i_{L_f}(\omega t) + i_{ac}(\omega t). \quad (3)$$

Differentiating both sides gives

$$\frac{1}{q_e^2} \frac{d^2 v_{S_n}(\omega t)}{d\omega t^2} + v_{S_n}(\omega t) - p_e \cdot \cos(\omega t + \varphi) = 0 \quad (4)$$

of which the general solution is given by

$$v_{S_n}(\omega t) = \xi_1 \cos(q_e \omega t) + \xi_2 \sin(q_e \omega t) - \xi_e q_e \cos(\omega t + \varphi) \quad (5)$$

where  $\xi_1$  and  $\xi_2$  are the coefficients related to the circuit parameters of the inverter.  $v_{S_n}$  is the switch voltage normalized by  $V_{dc}$ , and

$$\xi_e = \frac{q_e p_e}{1 - q_e^2} \quad (6)$$

$$q_e = \frac{1}{\omega \sqrt{L_e C_f}} \quad (7)$$

$$p_e = \frac{I_{ac} \cdot \omega L_e}{V_{dc}} \quad (8)$$

where  $p_e$  is defined as the load factor representing the output current limited by the load resistance. To find the expressions of  $\xi_1$ ,  $\xi_2$ , the initial conditions at  $\omega t = \pi$  can be used as follows.

The inherent zero-voltage turning-OFF condition of  $S_1$  is expressed as

$$v_{S_n}(\pi) = 0 \quad (9)$$

and the voltage-second balance of  $L_f$  is given by

$$\int_{\pi}^{2\pi} v_{S_n}(\omega t_1) d\omega t_1 = 2\pi. \quad (10)$$

Thus, the expressions of  $\xi_1$  and  $\xi_2$  are solved according to the two initial conditions, i.e.,

$$\begin{aligned} \xi_1 &= \frac{-q_e \xi_e}{1 - \cos(\pi q_e)} [(2q_e \sin(\varphi) + 2\pi/\xi_e) \sin(\pi q_e) \\ &\quad + \cos(\varphi) \cos(\pi q_e) (1 - 2\cos(\pi q_e)) + \cos(\varphi)] \\ \xi_2 &= \frac{q_e \xi_e}{1 - \cos(\pi q_e)} [(2q_e \sin(\varphi) + 2\pi/\xi_e) \cos(\pi q_e) \\ &\quad - \cos(\varphi) \sin(\pi q_e) (1 - 2\cos(\pi q_e))]. \end{aligned} \quad (11)$$

Hence, the normalized switch voltage  $v_{S_n}$  is expressed as a function of three variables ( $q_e$ ,  $p_e$ , and  $\varphi$ ), which are crucial for modeling the following criteria achieving load-independent operation. It is worth noting that (11) has fewer variables than that in [27], where the sets of four variables are calculated numerically.

The following section obtains a unique and unreported analytical parameter set of the proposed topology achieving load-independent soft-switching and constant voltage gain.

## B. Criteria of Load-Independent Operation

Since  $v_{S_n}$  acts a voltage source for the output. To analyze the output voltage with load variation,  $v_{S_n}$  can be separated into the active part ( $V_{R_n}$ ) and the reactive part ( $V_{X_n}$ ) by Fourier series as

$$V_{R_n} = \frac{1}{\pi} \int_{\pi}^{2\pi} v_{S_n}(\omega t) \sin(\omega t + \varphi) d\omega t \quad (12)$$

$$V_{X_n} = \frac{1}{\pi} \int_{\pi}^{2\pi} v_{S_n}(\omega t) \cos(\omega t + \varphi) d\omega t. \quad (13)$$

Deriving (12) and (13) with respect to the load factor  $p_e$  gives

$$\frac{dV_{R_n}}{dp_e} = \frac{q_e^4 (\cos(2\varphi) - 1) (\cos(\pi q_e) + 1)}{\pi (q_e^2 - 1)^2 (\cos(\pi q_e) - 1)} \quad (14)$$

$$\begin{aligned} \frac{dV_{X_n}}{dp_e} &= \frac{2q_e^4 \cos(\varphi) (\cos(\pi q_e) + 1)}{\pi (q_e^2 - 1)^2 (1 - \cos(\pi q_e))} \cdot \sin(\varphi) \\ &\quad + \left[ \frac{2q_e \sin(\pi q_e)}{\pi (q_e^2 - 1) (1 - \cos(\pi q_e))} + \frac{1}{2} \right] \cdot \frac{q_e^2}{q_e^2 - 1}. \end{aligned} \quad (15)$$

Let  $\varphi = \pi$  or 0 results

$$\left. \frac{dV_{R_n}}{dp_e} \right|_{\varphi=\pi,0} = 0. \quad (16)$$

The two particular values of  $\varphi$  lead to a same result. Therefore,  $V_{R_n}$  is irrelevant to load variation when  $\varphi = \pi$ . Then, (15) can be simplified as

$$\begin{aligned} \xi_4 &= \left. \frac{dV_{X_n}}{dp_e} \right|_{\varphi=\pi} \\ &= \left[ \frac{2q_e \sin(\pi q_e)}{\pi (q_e^2 - 1) (1 - \cos(\pi q_e))} + \frac{1}{2} \right] \cdot \frac{q_e^2}{q_e^2 - 1}. \end{aligned} \quad (17)$$

As  $q_e$  depends on circuit parameters,  $V_{X_n}$  is proportional to  $p_e$  defined in (8). The voltage gain  $G_v$  can be expressed as

$$G_v = \frac{V_{ac}}{V_{dc}} = 2 \cdot \left| j \cdot p_e \left( \xi_4 + \frac{L_x}{L_e} \right) + V_{R_n} \right| \quad (18)$$

where  $V_{ac}$  is the amplitude of  $v_{ac}$ . Thus, the reactive part of  $V_{S_n}$  can be compensated by fixed inductance given by

$$L_x = -\xi_4 L_e \quad (19)$$

Therefore, the compensated voltage-gain is given by

$$G_v = \frac{4q_e \sin(\pi q_e)}{(\cos(\pi q_e) - 1)(q_e^2 - 1)} \quad (20)$$

which is load-independent. The value of  $q_e$  achieving ZVS turn-ON can be obtained by solving  $v_{S_n}$  at  $\omega t = 2\pi$ , i.e.,

$$v_{S_n}(2\pi) = \frac{2\pi q_e \sin(\pi q_e)}{(1 - \cos(\pi q_e))} = 0 \quad (21)$$

which gives

$$q_e = 1. \quad (22)$$

Since the denominator of (15) cannot be zero, the limit method can be applied to evaluate the compensation coefficient, and

gives

$$\xi_4 = -\frac{1}{8} \quad (23)$$

which results in a compensated voltage gain irrelevant to load variation, i.e.,

$$G_v = \pi. \quad (24)$$

Meanwhile, the  $v_{S_n}$  can be simplified as

$$v_{S_n}(\omega t) = \sin(\omega t) \left[ p_e \left( \frac{3\pi}{4} - \frac{\omega t}{2} \right) - \pi \right] \quad (25)$$

thus, the switch voltage has no high-order harmonics at no-load condition ( $p_e = 0$ ).

As a result, (22) and (23) form a unique parameter set, which achieves load-independent ZVS turn-ON and maintains a constant voltage gain of  $\pi$  irrelevant to load resistance.

### C. Theoretical Parameters and Waveforms

Based on applicable specifications, the rated condition can be designed where the switches achieve both ZVS and zero-current-switching (ZCS) turn-ON. Thus, the inverters perform soft switching over the entire load range and prevent reverse conducting at the switches' OFF-state.

The value of load factor  $p_e$  achieving ZCS turn-ON can be obtained by solving

$$\left. \frac{dv_{S_n}}{d\omega t} \right|_{\omega t=2\pi} = -\frac{\pi}{4} \cdot p_e - \pi = 0 \quad (26)$$

which gives

$$p_{e_r} = -4. \quad (27)$$

This ZCS condition is involved to form an equation set for calculating the practical circuit parameters as

$$R_{ac_r} = \frac{V_{dc}^2 G_v^2}{2P_{out_r}} \quad (28)$$

$$R_{dc_r} = \frac{V_{dc}^2}{P_{out_r}} \quad (29)$$

$$L_f = \frac{R_{dc_r} G_v}{\omega} \quad (30)$$

$$L_x = 0.25L_f \quad (31)$$

$$C_f = \frac{1}{2L_f \cdot \omega^2} \quad (32)$$

where  $P_{out_r}$  is the inverter's rated output power.  $R_{dc_r}$  and  $R_{ac_r}$  are the rated input and load resistance, respectively. It can be seen that the circuit parameters have constant coefficients. Equations (28)–(32) provide a concise guideline to design practical systems.

The switch voltage at various loads can be calculated according to (25), where  $R_{ac}$  is inversely proportional to  $p_e$  due to the constant voltage gain. As in Fig. 2(a), the switch maintains ZVS as load varies. Meanwhile, it will conduct reversely during a particular interval when  $R_{ac}$  is greater than the rated value. For the MOSFET or GaN FET-based systems, the semiconductor

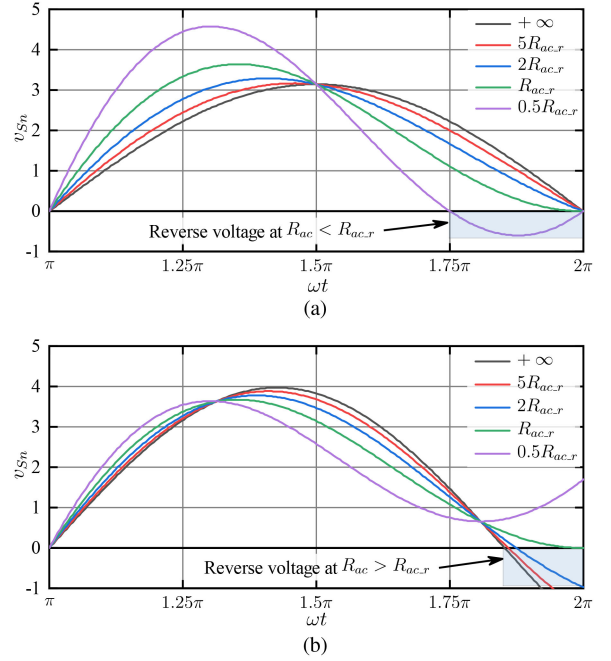


Fig. 2. Normalized switch voltage  $v_{S_n}$  versus load resistance  $R_{ac}$ . The two models perform noticeably different switching features. (a) Proposed model achieving load-independent ZVS (this article). The switches of proposed model achieve ideal ZVS turn-ON over the designed load range. (b) Conventional model in [27]. The switch in conventional model realizes ZVS turn-ON at the rated load but becomes non-ZVS at light loads.

devices have a limited reverse conduction voltage drop, which will cause additional losses and affect the voltage gain. Therefore, the inverter can prevent reverse conducting of the switches at OFF-state over the designed load range (i.e.,  $R_{ac}$  varies from the rated value to infinity). The peak voltage increases as  $R_{ac}$  decreases. The maximum  $v_{S_n}$  reaches 4.575 at  $0.5 \cdot R_{ac_r}$ , and 3.639 at  $R_{ac_r}$ . The reverse voltage usually causes unexpected conducting of practical semiconductor devices and affect the load-independent operation.

The switching behavior is quite different from that in [27]. As shown in Fig. 2(b) where the conventional model achieves ZVZCS turn-ON at rated load, but turns to non-ZVS at  $R_{ac} > R_{ac_r}$ . Using lower  $q_e$  can reduce switch voltage offsets of the conventional model and achieve better ZVS behavior. Nevertheless, the coupled inductor's current amplitude will substantially increase and cause unbearable losses.

In contrast, the proposed model achieves strict ZVS turn-ON over the entire load range without compromising the winding currents. Thus, the soft-switching performance is distinctly improved than before. The normalized  $i_{L_f}$  is given by

$$i_{L_f n}(\omega t) = \frac{i_{L_f}(\omega t)}{I_{dc_r}} = \frac{2}{p_{e_r} \cdot G_v} \cdot \begin{cases} \int_{\pi}^{\omega t} v_{S_n}(\omega t_1 + \pi) d\omega t_1 - \frac{\pi p_e}{4} + 1 & 0 < \omega t \leq \pi \\ -\int_{\pi}^{\omega t} v_{S_n}(\omega t_1) d\omega t_1 - \frac{\pi p_e}{4} + 1 & \pi < \omega t \leq 2\pi. \end{cases} \quad (33)$$

As illustrated in Fig. 3, the winding currents of  $L_{in}$  maintain an amplitude of  $0.5 \cdot I_{dc_r}$  at various loads, and it is sinusoidal at no-load condition. As the amplitude keeps relatively large over

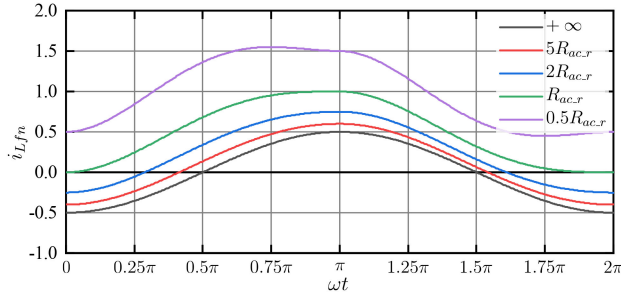


Fig. 3. Normalized winding current  $i_{Lfn}$  versus load resistance  $R_{ac}$ . The amplitude of  $i_{Lfn}$  is nearly constant as  $R_{ac}$  varies.

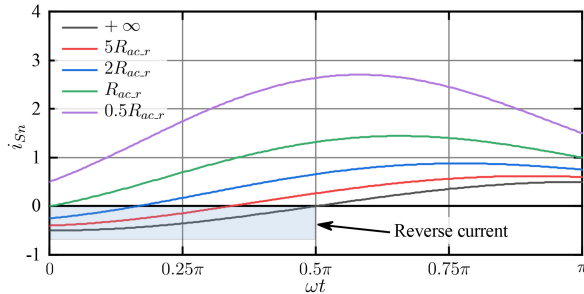


Fig. 4. Normalized switch current  $i_{Sn}$  versus load resistance  $R_{ac}$ . The amplitude increases as  $R_{ac}$  decreases, and the switches conduct reversely at the beginning of each ON-state period when  $R_{ac} > R_{ac,r}$ .

the entire load range, the coupled inductor occupies a substantial part in the total loss.

As illustrated in Fig. 4, the switch achieves both ZVS and ZCS turn-ON at the rated condition. Otherwise, the switch conducts reversely during a particular interval at the beginning of each ON-state period. The switch current  $i_S$  normalized by the rated input current  $I_{dc,r}$  is given by

$$i_{Sn}(\omega t) = \frac{i_S(\omega t)}{I_{dc,r}} = \begin{cases} \frac{2p_e}{p_{e,r} \cdot G_v} \cdot \sin(\omega t + \pi) + i_{Lfn}(\omega t) & 0 < \omega t \leq \pi \\ 0 & \pi < \omega t \leq 2\pi. \end{cases} \quad (34)$$

### III. LCC-S RESONANT CLASS E<sup>2</sup> TOPOLOGY FOR MHZ-WPT APPLICATIONS

This section proposes a push-pull Class E<sup>2</sup> topology as in Fig. 5 for MHz-WPT applications. The active Class E rectifier is implemented and has ZPA input irrelevant to load resistance. Besides, the LCC-S resonant tank realizes impedance conversion with zero phase shift. Eventually, the push-pull Class E<sup>2</sup> topology features load-independent ZVS operation and constant voltage gain.

#### A. Active Class E Rectifier With Coupled Inductor

As analyzed in Section II, the conventional inductor-coupled push-pull Class E topology has variable ZVS phases. Thus, it requires complex switching control when implemented as active rectification. Moreover, the rectifier's input phase varies with load resistance and causes complicated effects on the switching

behavior of the inverter on the transmitter [27]. On the other hand, the proposed load-independent topology has an invariable ZVS phase. Thus, the switching phase control can be simplified significantly for achieving soft-switching.

According to the time-domain modeling and Fourier series in Section II, the rectifier's input impedance  $R_{rec}$  is given by

$$R_{rec} = 0.5\pi^2 R_L \quad (35)$$

where  $R_{rec}$  maintains ZPA regardless of the load resistance  $R_L$  variation. The parameters for the active rectifiers are identical to that of the inverters.

As in Fig. 6, the synchronous switches maintain the ZVS turn-ON over the entire load range, including  $R_L < R_{dc,r}$ . Thus, it is worth noting that the active rectifiers perform better soft-switching behavior than the inverters. Besides, the ZCS turn-OFF is realized at  $R_L = R_{dc,r}$ . Similar to the inverters, the reverse voltage occurs during an interval at the beginning of each OFF-state period when  $R_L < R_{dc,r}$ . As in Fig. 7, the switches carry forward currents during an interval of each ON-state period. Thus, the load-independent rectification cannot be handled by diode-based rectifiers. Fig. 8 illustrates the currents through the output inductor  $L_{out}$ . It also features low-harmonic contents and large current amplitude at various loads.

#### B. LCC-S Resonant Tank

The LCC-S resonant tank can realize constant voltage gain and impedance conversion without phase shift. Thus, the inverter's load maintains resistive since the active rectifier has load-independent ZPA input. Thus, both the inverter and the active rectifier can maintain load-independent operation when connected through the LCC-S tank, of which parameters satisfy

$$\frac{1}{\omega^2} = L_1 C_1 = L_2 C_{2eq} = L_3 C_3 \quad (36)$$

where  $C_{2eq} = C_1 C_2 / (C_1 + C_2)$ , and thus, it comprises three resonant loops, which realizes a voltage gain expressed as

$$G_r = \frac{k_{23} \sqrt{L_2 L_3}}{L_1} \quad (37)$$

where  $k_{23}$  is the coupling coefficient between the transmitting coil and the receiving coil. As the push-pull Class E<sup>2</sup> topology is symmetric, the dc-dc voltage gain of the system is identical to the LCC-S tank, i.e.,  $G_{sys} = G_r$ . This equation is still valid when  $G_r \neq 1$  since both the inverter and the rectifier have constant voltage gains.

The proposed push-pull Class E<sup>2</sup> topology has a load-independent voltage gain, proportional to the coupling coefficient between the transmitting coil and the receiving coil.

#### C. Magnetic Structure of Coupled Inductors With Enhanced Self-Inductance

As analyzed above, the coupled inductors dominate the suppression of current ripples on dc sides. The equivalent self-inductance can be increased with winding turns or magnetic core volumes. However, the overall losses will substantially increase due to the large amplitudes of harmonic currents. Therefore, the

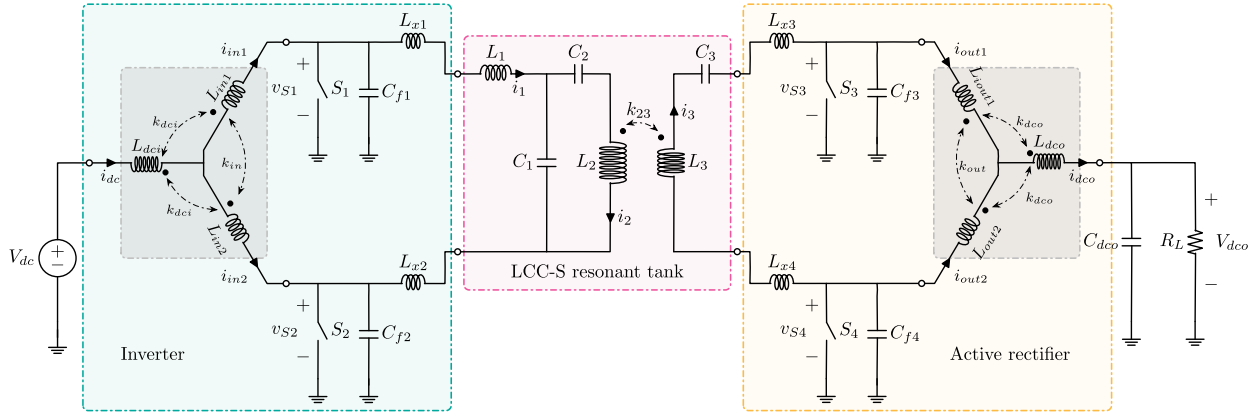


Fig. 5. Proposed push-pull Class E<sup>2</sup> topology with the three-winding magnetic structure for MHz-WPT applications. The additional inductance  $L_x$  can be merged with  $L_1$  and  $L_3$ , respectively.

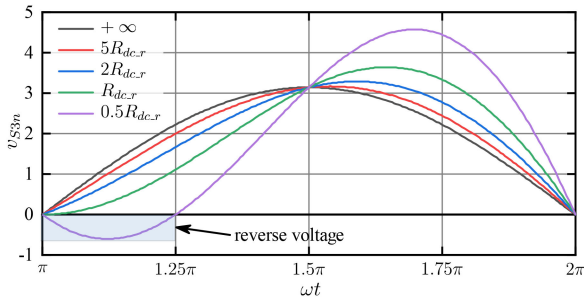


Fig. 6. Switch voltage  $v_{S3}$  of the active rectifier normalized by  $v_{dco}$ . The switches maintain soft-switching as  $R_L$  varies. Note that  $v_{S3}(\omega t) = v_{S4}(\omega t + \pi)$ .

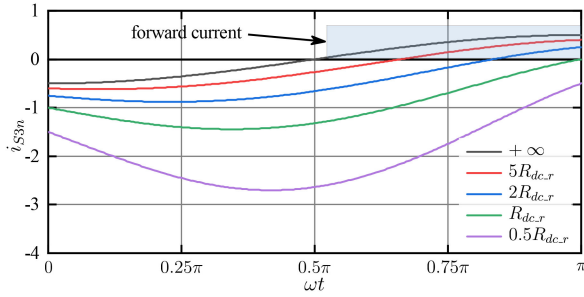


Fig. 7. Switch current  $i_{S3}$  of the active rectifier normalized by the rated  $I_{dco}$ . The forward current cannot be handled by diode-based rectifier.

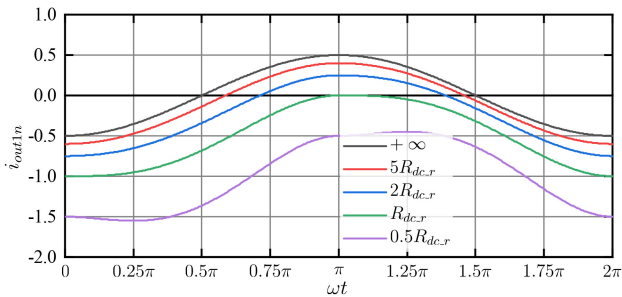


Fig. 8. Inductor current  $i_{out1}$  normalized by the rated  $I_{dco}$ . The current amplitude is nearly constant as  $R_{dc}$  varies.

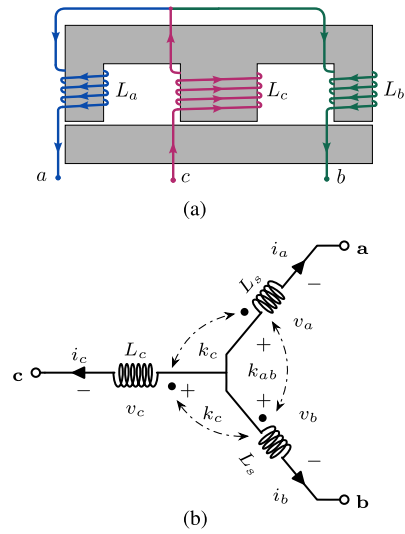


Fig. 9. (a) Proposed three-winding magnetic structure. (b) Equivalent circuit. Both the coupling coefficients are negative, i.e.,  $k_{ab}, k_c < 0$ .

conventional magnetic structures of two windings are difficult to implement sufficient self-inductance [27], [28].

As in Fig. 9, a three-winding magnetic structure is proposed to improve the coupled inductors. With the additional winding, the equivalent self-inductance can be increased individually without increasing the turns carrying large harmonic currents. As illustrated in Fig. 9(b), since the structure is symmetrical,  $L_a = L_b = L_s$ . The equivalent circuit gives

$$\begin{bmatrix} v_a \\ v_b \\ v_c \end{bmatrix} = \begin{bmatrix} L_s & k_{ab}L_s & k_c\sqrt{L_sL_c} \\ k_{ab}L_s & L_s & k_c\sqrt{L_sL_c} \\ k_c\sqrt{L_sL_c} & k_c\sqrt{L_sL_c} & L_c \end{bmatrix} \begin{bmatrix} \frac{di_a}{dt} \\ \frac{di_b}{dt} \\ \frac{di_c}{dt} \end{bmatrix}. \quad (38)$$

Rearranging the equations gives

$$\begin{bmatrix} v_{ca} \\ v_{cb} \end{bmatrix} = \begin{bmatrix} L_m + L_f & L_m \\ L_m & L_m + L_f \end{bmatrix} \begin{bmatrix} \frac{di_a}{dt} \\ \frac{di_b}{dt} \end{bmatrix} \quad (39)$$

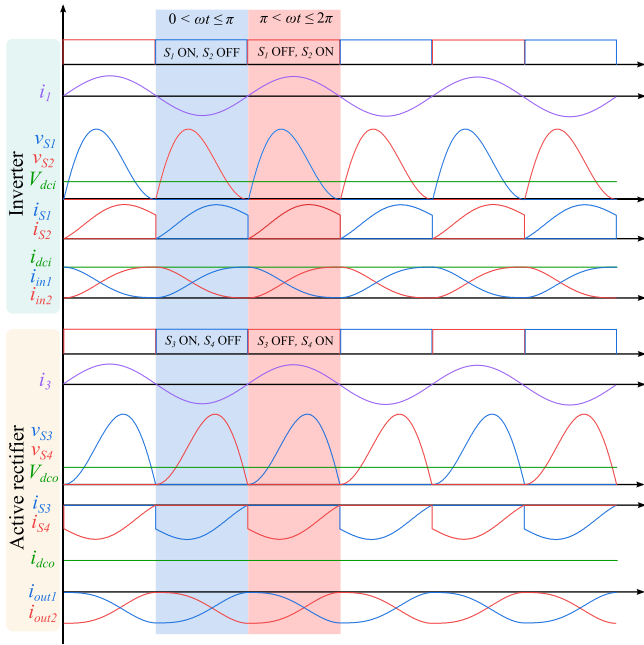


Fig. 10. Theoretical waveforms of the proposed  $LCC-S$  resonant Class  $E^2$  topology for MHz-WPT applications. It features low-harmonic contents and ideal soft switching. The theoretical input and output currents have zero ripples.

where  $L_m$  and  $L_f$  are the equivalent decoupled inductance in Fig. 1(b) given by

$$L_m = L_c + k_{ab}L_s - 2k_c\sqrt{L_sL_c} \quad (40)$$

$$L_f = (1 - k_{ab})L_s = \frac{L_{ab}}{2} \quad (41)$$

where  $L_{ab}$  is the inductance between terminals **a** and **b**. Thus, the equivalent self-inductance can be enhanced by increased  $L_c$  without increasing the turns of  $L_a$  and  $L_b$ . Meanwhile,  $L_f$  depends only on the coupled  $L_a$  and  $L_b$ , and is irrelevant to  $L_c$ . Eventually, the proposed magnetic structure provides a simplified method to implement the coupled inductors and obtain sufficient self-inductance for dramatically suppressing the harmonic currents on dc sides.

#### D. Theoretical Waveforms and Design Consideration

Fig. 10 illustrates the theoretical waveforms of the proposed Class  $E^2$  converter at the rated operation condition. The peak switch voltage will decrease as  $R_L$  increases. Thus, it reaches the maximum peak value of about  $3.6 V_{dc}$  at the rated load. Consequently, the processable power of the push-pull Class  $E^2$  topology is limited by the relatively high-voltage stress of the switches.

For practical applications, the expected voltage gain of the  $LCC-S$  network depends on specifications, which may also include limitations of size, weight, etc. According to (37), various combinations of the parameters are available to realize an identical voltage gain and will lead to different efficiency. To improve the efficiency, the loaded quality factor can be appropriately

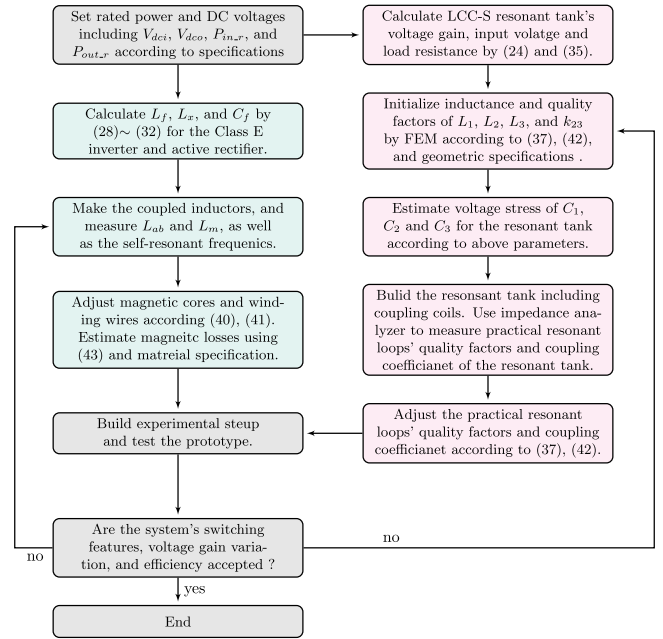


Fig. 11. Process of designing a practical MHz-WPT system based on the proposed load-independent Class  $E^2$  converter with coupled inductors.

adjusted according to the approximate constraint given by

$$Q_L = \frac{\omega L_3}{R_{ac}} \approx \frac{1}{k_{23}\sqrt{\frac{Q_2}{Q_3} + G_r k_{23}\frac{Q_2}{Q_1}\sqrt{\frac{L_2}{L_3}}}} \quad (42)$$

which is simplified by assuming the quality factors of the resonant loops are much larger than  $\frac{1}{k_{23}}$ . This equation provides a concise way to arrange the coupling coils, of which parameters depend on the sizes and the numbers of turns.

As analyzed above, both  $i_{in}$  and  $i_{out}$  have fixed current amplitudes of about  $0.5 I_{dc}$ . Thus, the magnetic core losses are dominated by the side columns. According to (30), the amplitude of magnetic flux density passing the side columns can be estimated by

$$B_m \approx \frac{1 + k_{ab}}{1 - k_{ab}} \cdot \frac{V_{dc}}{4fN_fA_e} \quad (43)$$

where  $f$  is the operation frequency, and  $A_e$  is the effective cross-sectional area of the side columns, as illustrated in Fig. 9(a). The acceptable  $B_m$  depends on the loss characteristics of the magnetic cores, of which manufacturers usually provide loss density curves versus  $B_m$ . Then, the number of turns for  $L_f$  can be properly selected according to (41) and (43).

As shown in Fig. 5, a practical MHz-WPT system comprises the inverter, active rectifier, and  $LCC-S$  resonant tank. To summarize the utilization of the above equations, a brief design process is illustrated in Fig. 11, which integrates the modeling presented before.

#### IV. EXAMPLE DESIGN AND IMPLEMENTATION

This section is to verify the proposed methodology. A 6.78-MHz prototype is designed and implemented to realize unity voltage-conversion ratio by using the topology in Fig. 5.

TABLE I  
 PARAMETERS OF INVERTER AND ACTIVE RECTIFIER

$L_{ab}$	1.133 $\mu\text{H}$
$L_x$	142 nH
$C_f$	486 pF
Operation frequency	6.78 MHz
GaN FET ( $S_{1\sim 4}$ )	GS66508T

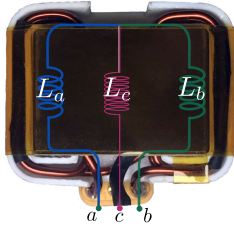


Fig. 12. Coupled inductor using EEQ25/11 magnetic core of Fair-Rite 67 material. A 3-D-printed skeleton is applied to fabricate the three windings.

### A. Inverter and Active Rectifier

According to the design process illustrated in Fig. 11, the prototype's parameters in Table I are calculated by (29)–(32) and (41) with expected rated power and dc voltages. The inductance values have a straightforward rule as  $L_f = 0.5L_{ab} = 4L_x$ , which simplifies the design of the coupled inductors and the additional inductor.

For simplifying the design and comparison, the inverter and the rectifier have the same rated power and dc voltages, i.e.,  $P_{in,r} = P_{out,r} = 300 \text{ W}$ ,  $V_{dc,i} = V_{dc,o} = 48 \text{ V}$  in this design example. As analyzed above, the load-independent operation is conducive to preventing reverse conducting of GaN FETs at OFF-state. The switches are expected to achieve ZVS turn-ON over the designed load range.

As in Fig. 12,  $L_{in}$  and  $L_{out}$  are implemented by the same magnetic structure illustrated in Fig. 9. Considering  $i_c$  are dominated by direct current, the winding of  $L_c$  is made by solid copper wires of 0.6 mm diameter. In addition, since Litz wires are no longer effective for reducing winding losses at 6.78 MHz, solid wires of 1.2 mm diameter are applied for  $L_a$  and  $L_b$ , of which current amplitude is nearly constant at various loads. The measured  $L_{ab}$  is approximately 1  $\mu\text{H}$ , and the self-inductance (i.e.,  $L_{ac}$ ,  $L_{bc}$ ) is about 5  $\mu\text{H}$ , which provides enough reactance to suppress the high-frequency current ripples.

A phase detector is implemented to detect and synchronize the driving phases of the synchronous switches. Since the input impedance of the active rectifier maintains ZPA regardless of load variation, the induced voltage of the receiving coil has a constant phase shift from the expected driving pulses. In this example, an auxiliary coil is applied to sense the induced voltage phase of the receiving coil [19]. As in Fig. 13, the auxiliary coil comprises three turns wound in different directions. By adjusting the pitches, the auxiliary coil is decoupled with the receiving coil but maintains a small mutual inductance with the transmitting coil. Then, the auxiliary coil can detect the induced voltage phase of the receiving coil, and avoid interference from the power stage. Two stages of monostable multivibrators are implemented to adjust the driving pulses' delay and duty cycles. As a result,

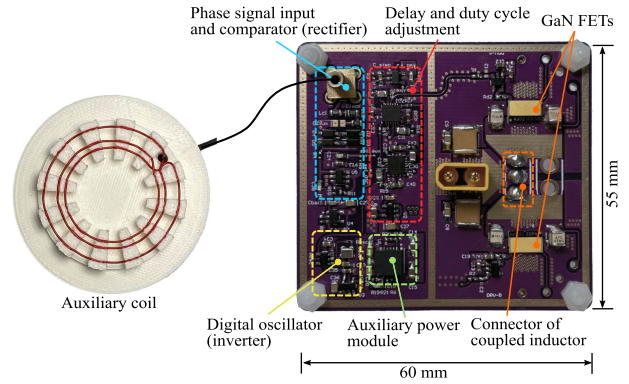


Fig. 13. Printed circuit board for the inverter and the active rectifier. The auxiliary coil provides an initial signal of phase reference for the synchronous switches of the rectifier. For the inverter, the driving pulses are generated by the digital oscillator on the PCB.

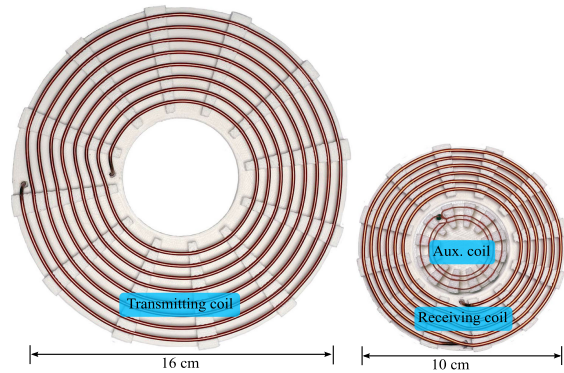


Fig. 14. Coupling coils. The auxiliary coil is fixed on the receiving coil. It is decoupled with the receiving coil but maintains a small mutual inductance with the transmitting coil.

the synchronous switches can be driven appropriately using the auxiliary coil and signal conditioning circuit.

The printed circuit board (PCB) in Fig. 13 is also implemented for the inverter. As phase synchronization is not necessary for inversion, a digital oscillator on PCB is applied to generate the driving pulses directly. Besides, a dc–dc module (TPSM5601) is used to power the gate drivers (1EDN7512) and the signal conditioning circuit.

### B. LCC-S Compensated Coupling Coils

The geometry of coupling coils depends on particular applications. Generally, increasing coil size will result in lower self-resonant frequency (SRF). To compensate the coils properly and ensure the expected operation of the resonant tank, the SRF of the coupling coils should be much higher than 6.78 MHz. As in Fig. 14, the coupling coils are made by solid copper wires. The outer diameters of the transmitting coil and the receiving coil are 16 and 10 cm, respectively, and the pitches were adjusted according to full-wave finite element method (FEM) simulation. Eventually, the measured SRFs of the coupling coils are higher than 30 MHz, which satisfies the requirement of this example.

Table II illustrates the measured parameters of the resonant tank. The quality factors of  $L_1$ ,  $L_2$ , and  $L_3$  are 300, 371, and 446 at 6.78 MHz, respectively. To realize 1:1 voltage gain,

TABLE II  
PARAMETERS OF LCC-S RESONANT TANK

Transmitting coil	Outer diameter	16 cm
	Number of turns	8
	Pitch	6 mm
	Wire diameter	1.8 mm
	Inductance $L_2$	10.33 $\mu$ H
Receiving coil	Outer diameter	10 cm
	Number of turns	6
	Pitch	4 mm
	Wire diameter	1.7 mm
Compensation network	Inductance $L_3$	3.94 $\mu$ H
	$L_1$	1.82 $\mu$ H
	$C_1$	303.3 pF
	$C_2$	62.8 pF
	$C_3$	139.9 pF

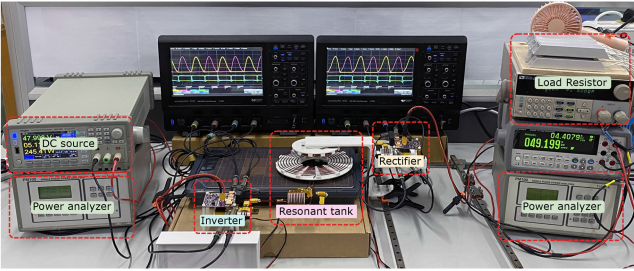


Fig. 15. Schematic diagram of the experimental setup.

the expected coupling coefficient meets  $k_{23} = \frac{L_1}{\sqrt{L_2 L_3}} = 0.286$ , corresponding to a transfer distance of about 2.5 cm. Thus, according to (42), the peak efficiency of the resonant tank is up to 97.8% at  $Q_L = 2.952$ , at which the output power is 200 W.

## V. EXPERIMENTAL MEASUREMENT AND ANALYSIS

Fig. 15 shows the experimental setup of which the prototype comprises the inverter, the LCC-S resonant tank, and the active rectifier. The coupled inductors are connected under the PCBs. To limit the temperature rise of the switches and the inductors, single-layer copper foils are implemented as heat sinks of the GaN FETs with air cooling fans.

As the receiver has no external power source, the rectifier operates at passive rectification mode for milliseconds until the auxiliary power module is settled. The reverse-conducting threshold of the GaN FETs is about 2 V at  $V_{gs} = 0$ . Thus, long-time passive rectification will degrade efficiency substantially due to the relatively high voltage drop of switches.

All the following measurements are obtained without adjusting circuit components. The relative position of the coupling coils is fixed with a measured  $k_{23} = 0.29$ . The operation frequency and the duty cycle are set at 6.78 MHz and 50%, respectively. Besides, the output voltage is fixed to 48 V by slightly adjusting the input voltage.

Fig. 16 illustrates the measured  $V_{ds}$  and  $V_{gs}$  of the inverter at various loads. The switches maintain ZVS turn-ON when the input power  $P_{in}$  is below 300 W, and the peak  $V_{ds}$  increases with the power level. The  $V_{ds}$  curves cross at  $\omega \approx 1.5\pi$  during the OFF interval of each period. When  $P_{in} = 294.4$  W, the voltage derivative of  $V_{ds}$  at turn-ON instant is near zero. Thus, the

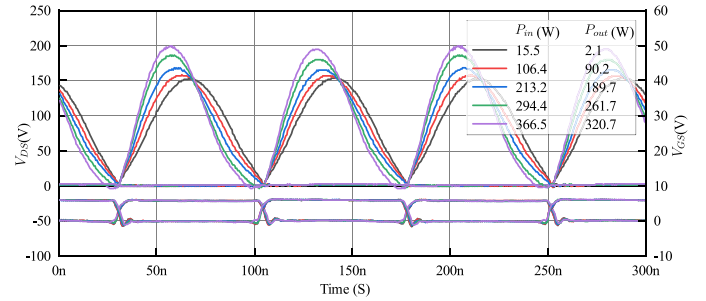


Fig. 16. Measured  $V_{ds}$  and  $V_{gs}$  of the inverter. The ZVS is maintained over the designed load range, i.e.,  $P_{in} \leq 300$  W.

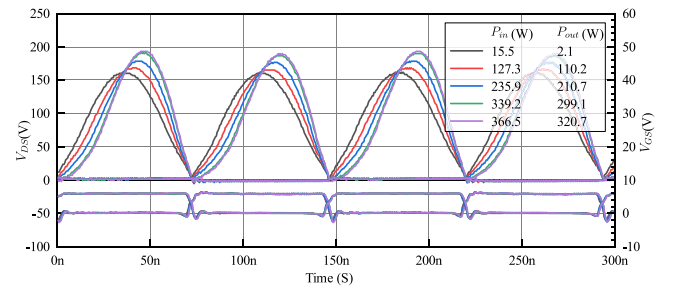


Fig. 17. Measured  $V_{ds}$  and  $V_{gs}$  of the active rectifier. The ZVS is maintained over the entire load range, including  $P_{out} > 300$  W.

switches achieve both ZVS and ZCS turn-ON when  $P_{in}$  is close to the rated input power. Otherwise, when the input power exceeds 300 W, the switches will conduct reversely during a short interval before the rising edge of each turn-ON pulse. Therefore, the measured  $V_{ds}$  matches well with the calculated curves illustrated in Fig. 2.

As illustrated in Fig. 17, the active rectifier acts as a time-reversal inverter. The measured  $V_{ds}$  curves also cross at  $\omega \approx 1.5\pi$ , and show a nearly mirrored image of the inverter. The switches maintain ZVS turn-ON over the entire load range, including  $P_{out} > 300$  W. At  $P_{out} = 320.7$  W as the output power exceeds the rated value,  $V_{ds}$  does not rise immediately, which means the synchronous switches conduct reversely during a short interval after the rising edge of each turn-OFF pulse. However, the OFF-state switches can commute naturally and realize ZCS turn-OFF. Thus, the active rectifier performs better soft switching behavior than the inverter. Comparing the measured  $V_{ds}$  to that calculated in Fig. 6, the rectifier shows consistent switching behavior with the analysis before.

Fig. 18(a) shows measured  $i_{in1}$  and  $i_{out1}$  by a current probe of 120 MHz bandwidth. According to the modeling before, the theoretical  $i_{Lf}$  maintains an amplitude of  $0.5I_{dc,r}$ , where  $I_{dc,r} = 6.25$  A for this design example. The measured currents are approximately sinusoidal, which match well with the calculated ones. It can be seen that the measured current amplitudes are slightly higher than the theoretical values illustrated in Figs. 3 and 8. In order to compare the measured currents with the analytical modeling, Fig. 18(b) illustrates the corresponding fast-Fourier transform (FFT) results. The fundamental current amplitudes maintain about 3.6 A at various loads. Besides, the dc offset increases with power level, and reaches about 3.2 A at the

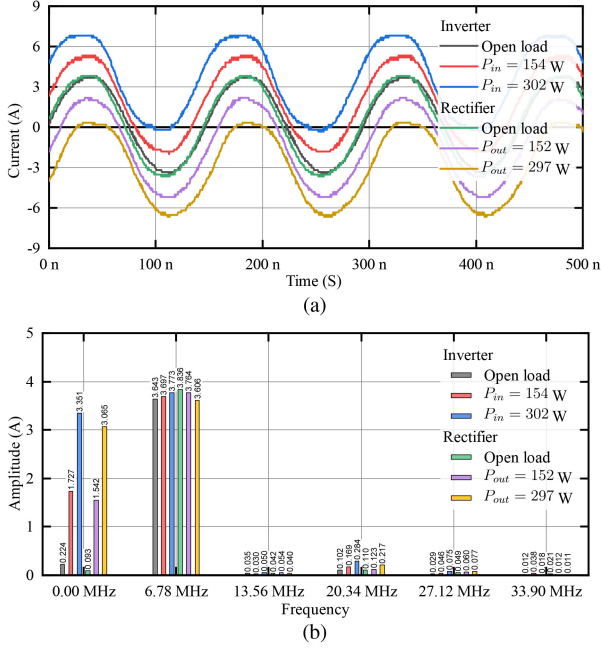


Fig. 18. (a) Measured  $i_{in1}$  and  $i_{out1}$ . (b) FFT results. The push-pull Class E<sup>2</sup> topology features low-harmonic contents.

rated input/output power of 300 W. Considering the measured  $L_{ab}$  of about  $1 \mu\text{H}$  is slightly lower than the calculated inductance of  $1.133 \mu\text{H}$  in Table I, the winding currents performed expected features as analysis before. The spectrum is dominated by dc and fundamental frequency at 6.78 MHz of which amplitudes are nearly constant at various loads. The even-order harmonics are well suppressed as the amplitudes are negligible at 13.56 and 27.12 MHz. The three-winding magnetic structure performs better suppression performance than that in [27]. As in Fig. 9, the fundamental current circulates between  $L_a$  and  $L_b$ , and thus the measured harmonic distribution implies that the ripples of  $i_{dci}$  and  $i_{dco}$  are negligible. It validates that the proposed magnetic structure is conducive to suppressing current noises on the dc sides and performs better filtering features than the conventional push-pull Class E topology with coupled inductors.

Fig. 19 shows the steady-state thermal images. The overall temperature rise increases with processed power and is dominated by the switches and the coupled inductors. The temperature rise of the GaN FETs shows slightly unbalanced due to the relatively high temperature of the auxiliary power module TPSM5601. According to the circuit modeling above, the fundamental current has to maintain a relatively high amplitude to realize load-independent ZVS operation. Consequently, the Class E<sup>2</sup> converter has relatively high no-load losses, which will reduce the efficiency at light-load conditions.

Fig. 20 shows the efficiencies and the voltage gain of the prototype. The output dc voltage is fixed at 48 V by slightly adjusting the input power supply. The maximum system efficiency is 89.3% at 210-W output power, and it will reach 90.4% when excluding the driving losses of which the measured value is 1.42 W for each side. Since the inverter and the active rectifier use the same parameters and present nearly symmetrical waveforms,

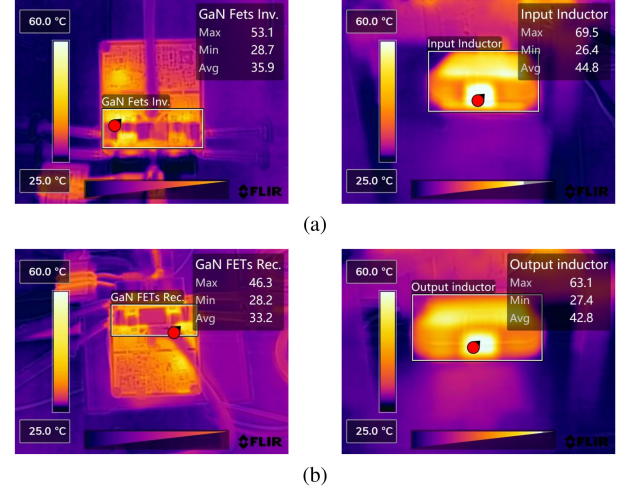


Fig. 19. Thermal images of the prototype with 339 W input to 300 W output. (a) Inverter. (b) Active rectifier. The total losses are dominated by the switches and the coupled inductors.

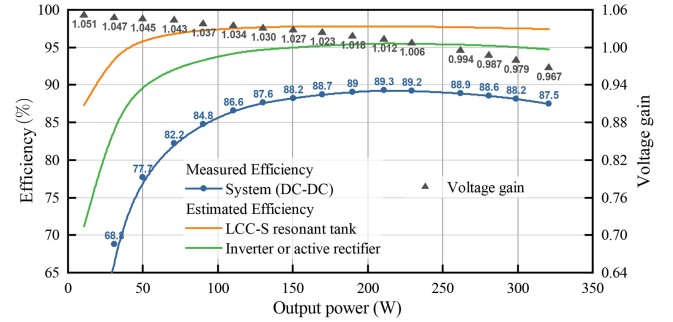


Fig. 20. Efficiencies and voltage gain.

it is reasonable to assume that they have approximately equal efficiencies. Therefore, the efficiencies of the LCC-S resonant tank, the inverter, and the rectifier are also estimated in the figure. The maximum efficiency of the inverter (or the active rectifier) reaches 95.6%.

As modeling in Section II, the load-independent operation achieves constant voltage gain for both the inverter and active rectifier. Considering the power losses of practical systems, the measured voltage gain will decrease with output power. As illustrated in the figure, the measured voltage gain, i.e.,  $V_{dco}/V_{dci}$ , decreases slowly with the output power. It varies between +5/-4% over the entire load range. Hence, the prototype presents the expected load-independent operation of the push-pull Class E<sup>2</sup> topology with coupled inductors.

## VI. CONCLUSION

This article models and implements the inductor-coupled push-pull Class E topology achieving the load-independent operation, which performs distantly improved soft-switching performance, load regulation, harmonics suppression, and efficiency compared with the previous work [27]. A prototype is built and the proposed methodology is verified. The following conclusions can be made.

- 1) The modeling validates that the inductor-coupled Class E topology is able to achieve the load-independent operation, which performs strict ZVS and constant voltage gain regardless of load variation. The measured switch voltages match well with the theoretical calculated.
- 2) The active rectifier performs better soft-switching than the inverter. The analysis and measurement show that the rectifier achieves load-independent ZVS even when the output power exceeds the rated value.
- 3) The proposed three-winding magnetic structure realizes much higher equivalent self-inductance, which dominates the suppression of harmonic currents. The measured FFT results indicate that the prototype features negligible current ripples on dc sides.
- 4) The proposed *LCC-S* resonant Class E<sup>2</sup> converter features negligible harmonic contents and maintains soft switching with a constant voltage conversion ratio over the entire load range. The prototype realizes load-independent ZVS and maintains an approximately constant voltage gain with +5/−4% variance from no-load to 320-W output power. The measured dc–dc efficiency reaches 89.3% at 210 W.

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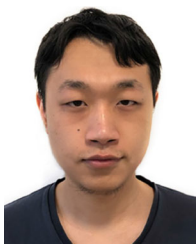
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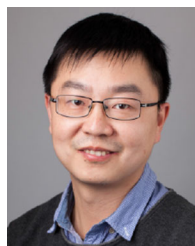
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