

# A Postprocessing-Technique-Based Switching Loss Estimation Method for GaN Devices

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**Abstract**—Gallium nitride (GaN) high electron mobility transistor (HEMT) is a promising candidate for the high-density power converter applications. Due to the low switching loss, the GaN HEMT may lead to a new horizon in the applications, such as fast charger, wireless charging, and 5G power amplifier. However, the fast switching speed of GaN HEMT makes it extremely sensitive to parasitic parameters, especially for the low-voltage GaN devices with small packages. Conventional switching loss estimation methods, which are normally based on the calibration fixture or modification of test circuit, cannot be effectively utilized for GaN HEMT. This article presents an accurate switching loss estimation method using the postprocessing technique. The proposed switching loss technique consists of three parts of signal processing, which are the wavelet denoising process to minimize the influence of background noise, the voltage–current ( $V - I$ ) alignment process to reduce errors from probe propagation delay, and the linear interpolation process for further improvement of alignment accuracy. In addition, a modified SPICE model is proposed to investigate the influence of parasitic parameters on switching losses by the circuit simulation. Based on the theoretical switching characteristics, the switching loss estimation method is finally validated experimentally on a commercial 40-V/10-A GaN HEMT in a double-pulse test.

**Index Terms**—Double-pulse test, gallium nitride (GaN), signal processing, switching characterization, switching loss estimation.

## I. INTRODUCTION

GALLIUM nitride (GaN) high electron mobility transistor (HEMT) is a promising candidate for the high-density power converter applications, such as fast charger, wireless charging, and 5G power amplifier [1]–[3]. Due to the advantages

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of wide bandgap (WBG) material, it has lower ON-resistance, lower parasitic capacitance, higher junction temperature, and higher switching speed compared with traditional silicon (Si) devices. As a lateral device, the GaN HEMT is fabricated on AlGaN/GaN heterojunction interface, which generates the two-dimensional electron gas (2DEG). It eliminates the p-n junction and leads to the zero reverse recovery charge ( $Q_{rr}$ ). Thus, it allows the GaN HEMT to operate in the MHz-frequency power converter with a high energy conversion efficiency [4]–[7]. When the switching frequency enters the MHz range, the switching loss of GaN HEMT becomes significant.

However, most of the device datasheets only provide limited information on the switching loss. Furthermore, the loss is strongly dependent on the circuit parameters. As a consequence, the switching loss characterization becomes crucial for the GaN-based converter design. So far, early works on the switching loss characteristics of GaN HEMT can be mainly classified into model-based [8]–[13] and measurement-based approaches [14]–[18].

The model-based approach requires accurate derivation of the behavioral or analytical models [9], [10], which can be a time-consuming and cumbersome process. Therefore, the measurement-based approach is usually preferred. This approach adopts two methods, the calorimetric and the double-pulse test (DPT) methods. The calorimetric method uses a calorimeter (such as an air flow meter [14]) to obtain the switching loss through capturing the temperature change, but it is challenging to obtain the breakdown of various loss mechanisms [18]. The DPT method acquires the time-domain voltage and current switching waveforms and performs the integration over time. Although there are drawbacks of DPT [19], such as the requirements of high-precision probes and an exact voltage–current ( $V - I$ ) alignment method, it is still considered as the most feasible method with reasonable accuracy [20]. By applying postprocessing technique, the influences from the probes and voltage–current misalignment can be minimized to improve the accuracy [21], [22].

Conventionally, probe manufacturers provide dedicated calibration fixtures for measuring instruments but only compatible with selected probes [23]. For example, the compensation output of an oscilloscope was used as the reference signal but it is only compatible to a limited selection of voltage or current probes [24]. Alignment of  $V - I$  relationship of the switching waveforms based on the  $Ldi/dt$  voltage drop resulted from the stray inductance has been reported but it is only effective

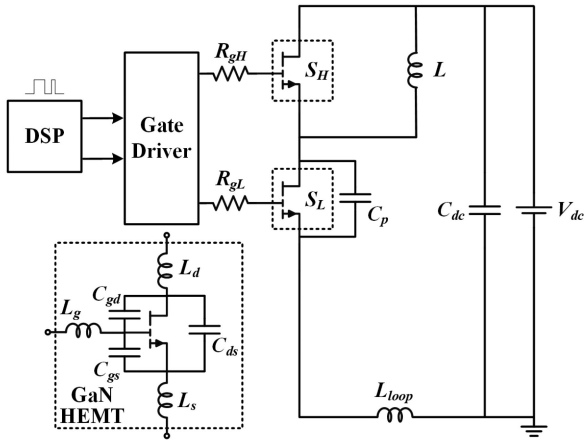


Fig. 1. Schematic of DPT circuit.

for sufficiently high load current [25]. Another reported work in [26] removes the freewheeling diode and load inductor and replaced them with a precision 100- $\Omega$  resistor with low parasitic inductance, which results in no phase shift between voltage and current waveforms alignment but the use external auxiliary circuit can be rather complex.

To overcome the aforementioned constraints, this article proposes a novel postprocessing method to estimate the switching loss accurately based on the DPT method. A postprocessing deskew technique for GaN devices is adopted. It consists of three parts, the wavelet denoising process, the  $V - I$  alignment process, and the linear interpolation process. First, to effectively reduce the error from background noise, the wavelet shrinkage is included in the wavelet denoising process. Then, based on the theoretical switching characteristics of GaN devices, the reference points of the switching waveform can be established so that the propagation delay of the probe could be eliminated during the  $V - I$  alignment process. Finally, the linear interpolation process is applied to improve the alignment resolution.

The rest of this article is organized as follows. Section II provides an overview of the theoretical loss mechanism of GaN HEMT. Section III develops a modified SPICE model of a commercial GaN HEMT EPC2014 to include the effects of parasitic elements in the switching loss. Next, a switching loss estimation method based on the proposed postprocessing technique is described in details in Section IV. Section V elaborates the DPT experiment setup for the GaN HEMT and Section VI validates the proposed loss estimating method experimentally. Finally, Section VII concludes this article.

## II. LOSS MECHANISM OF GaN DEVICES

For the DPT circuit shown in Fig. 1, the loss mechanism of low-side GaN HEMT ( $S_L$ ) will be analyzed to provide the *a priori* knowledge for the proposed method.  $L_{loop}$  is the power loop inductance of printed circuit board (PCB).  $C_p$  is the PCB parasitic capacitance.  $R_{gx}$  is external gate resistance.  $L$  is the load inductance.  $C_{dc}$  is dc-link capacitance, and  $V_{dc}$  is dc-bus voltage.  $C_{gs}$ ,  $C_{gd}$ , and  $C_{ds}$  are the gate-source, gate-drain, and drain-source junction capacitances, respectively.  $L_g$ ,  $L_d$ ,

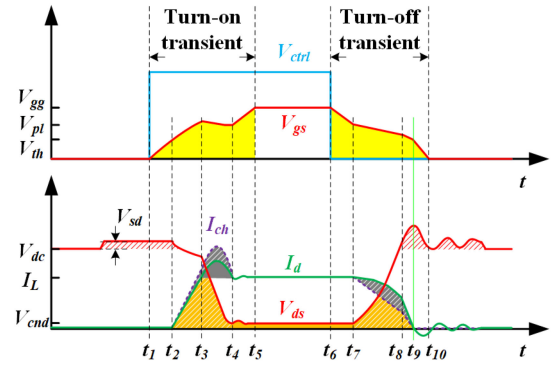


Fig. 2. Switching waveforms.

and  $L_s$  are the gate, drain, and source packaging inductances, respectively. The losses can be further classified into the turn-ON loss, the turn-OFF loss and the conduction loss, which will be discussed in details in the following sections. Definitions of these losses are based on switching waveforms illustrated in Fig. 2, where  $V_{ctrl}$ ,  $V_{gs}$ ,  $V_{ds}$ ,  $I_{ch}$ ,  $I_d$  are the control signal, gate-source voltage, drain-source voltage, 2DEG channel current, and drain current, respectively. In addition, parameter symbols of  $S_H$  are marked with an additional subscript “1”. It should be mentioned that due to the zero  $Q_{rr}$  property of GaN devices, there will be no reverse recovery loss  $E_{rr}$ . Although some literatures [11] have reported there is another type of loss referred to the output capacitance hysteresis, which is not considered in this article, because it only appears under very high-frequency (such as radio-frequency) application. Moreover, the impact of ringing on switching losses can be neglected, as it only contributes an insignificant part of the total loss [27].

### A. Loss During Turn-ON Transition

According to switching waveforms given in Fig. 2, the turn-ON transition of  $S_L$  can be subdivided into four stages, the turn-ON delay, the current commutation, the voltage commutation, and the current ringing. Each of these stages will be described in details to present the loss mechanism in the turn-ON transition. It includes the turn-ON  $V - I$  overlapping loss  $E_{vi,on}$  (orange zone with white stripe), the capacitor self-charging/discharging loss  $E_{oss}$  (gray zone with white stripe), the capacitor charging loss  $E_{qoss}$  (gray zone), the reverse conduction loss  $E_{sd}$  (white zone with red stripe), the forward conduction loss  $E_{cond}$  (orange zone), and the gate drive loss  $E_{gdv}$  (yellow zone).  $E_{cond}$  and  $E_{gdv}$  are not considered here for ease of analysis.

1) *Stage 1 (Turn-ON Delay,  $t_1 \sim t_2$ ):* When the control signal  $V_{ctrl}$  becomes high, the gate driver charges the input capacitance  $C_{iss}$  with a gate current  $I_g$  and  $V_{gs}$  rises. Since the value of  $V_{gs}$  is below the gate threshold voltage  $V_{th}$ ,  $S_L$  works in the cut-off region and its 2DEG channel is not conducted. Meanwhile,  $S_H$  is still in the reverse conduction, and therefore, losses in this stage only contribute to  $E_{sd}$ .

2) *Stage 2 (Current Commutation,  $t_2 \sim t_3$ ):*  $V_{gs}$  increases exponentially and exceeds  $V_{th}$ , and  $S_L$  enters the saturation region. The 2DEG channel starts to conduct, and the inductor current is transferred from  $S_H$  to  $S_L$ . Thus,  $E_{sd}$  still exists in

this stage and  $I_d$  is given by

$$I_d = g_m(V_{gs} - V_{th}) \quad (1)$$

where the transconductance  $g_m$  can be extracted from the transfer characteristics. The rise of  $I_d$  leads to a voltage drop across  $S_L$  due to the effect of stray inductance ( $L_{stray} = L_{d1} + L_{s1} + L_d + L_s + L_{loop}$ ) and  $V_{ds}$  is given by

$$V_{ds} = V_{dc} + V_{sd} - L_{stray} \frac{dI_d}{dt} \quad (2)$$

where  $V_{sd}$  is the reverse conduction voltage drop.  $C_{ds}$  discharges and  $V_{ds}$  falls, meanwhile,  $C_{gs}$  is charged by the current  $I_g$  continuously. Therefore, the increase of  $V_{gs}$  and decrease of  $V_{ds}$  lead to an increase of  $V_{gd}$ , thus charging  $C_{gd}$ . Since the 2DEG channel already conducts in this stage, the discharging current of  $C_{ds}$  and the charging current of  $C_{gd}$  will flow through the channel and the channel current  $I_{ch}$  is given by

$$I_{ch} = I_d + C_{oss} \frac{dV_{ds}}{dt} + C_{gd} \frac{dV_{gd}}{dt} \quad (3)$$

where  $C_{oss}$  is the output capacitance and given by  $C_{oss} = C_{gd} + C_{ds}$ . Considering  $C_{gd} \ll C_{ds}$  and  $dV_{gs}/dt \ll dV_{ds}/dt$ , the last term of (3) can be neglected. The first and second terms contribute to  $E_{vi,on}$  and  $E_{oss}$ , respectively. Apart from  $C_{oss}$ ,  $C_p$  also contributes to  $E_{oss}$ , which is voltage-independent [28] and is introduced by adjacent coppers [29], which can be extracted by the finite-element method (FEM) simulation. Since  $L_d$  and  $L_s$  of  $S_L$  are very small,  $C_{oss}$  and  $C_p$  can be combined into the equivalent capacitance  $C_{eq}$ . Thus, (3) can be simplified to

$$I_{ch} = I_d + C_{eq} \frac{dV_{ds}}{dt}. \quad (4)$$

In summary, this stage contributes to losses  $E_{sd}$ ,  $E_{vi,on}$ , and  $E_{oss}$ .

3) *Stage 3 (Voltage Commutation,  $t_3 \sim t_4$ ):* When  $I_d$  reaches the load current  $I_L$  and continues to rise, it leads to the further fall of  $V_{ds}$ .  $V_{gs}$  equals Miller plateau voltage  $V_{pl}$ . Similar to the analysis of stage 2, the voltage drop across  $C_{gs}$  reaches a dynamic balance while  $C_{gd}$  shunts the majority of charging current and increases nonlinearly, and  $V_{ds}$  drops, which is given by

$$V_{ds} = V_{dc} - V_{ds1} - L_{stray} \frac{dI_d}{dt} \quad (5)$$

where  $V_{ds1}$  is the drain–source voltage of  $S_H$ . Since the 2DEG channel of  $S_L$  is already conducting, the rest of self-discharging current of  $C_{eq}$  will flow into the 2DEG channel of  $S_L$ , which contributes to  $E_{oss}$ . Meanwhile, as the 2DEG channel (reverse biased) of  $S_H$  is already cut off, the charging current of  $C_{oss1}$  will flow through  $S_L$  and generate a current bump, which is the source of  $E_{qoss}$ , as shown in Fig. 3(a). Since this current bump is mixed with  $I_d$ ,  $E_{qoss}$  is considered as part of  $E_{vi,on}$  to simplify the analysis and  $I_d$  is now given by

$$I_d = I_L + C_{oss1} \frac{dV_{ds1}}{dt} \quad (6)$$

and  $I_{ch}$  is the same as in (4). To summarize, this stage contributes to losses  $E_{vi,on}$  and  $E_{oss}$ .

It should be noted that the voltage commutation can be eliminated in stage 3 under the extreme condition of low dc-bus

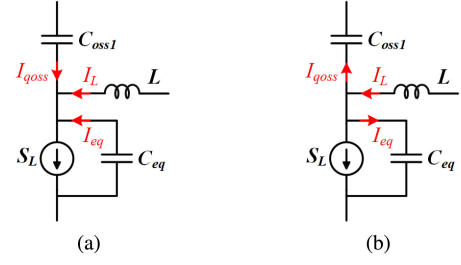


Fig. 3. Simplified circuit diagrams of current sharing in voltage commutation. (a) Turn-ON transition. (b) Turn-OFF transition.

voltage and large stray inductance. It can be explained by (2) that a large  $L_{stray}$  leads to last term comparable to  $V_{dc}$  and  $V_{ds}$  may drop to zero in stage 2.  $E_{vi,on}$  disappears in stage 3 since there is no  $V - I$  overlapping. Also,  $E_{oss}$  will be fully dissipated in stage 2. The derivation of turn-ON energy under extreme situation is given in the Appendix. Under the extreme condition, for example,  $V_{dc} = 18$  V and  $L_{stray} = 12.25$  nH,  $di/dt = 2.51$  kA/ $\mu$ s, the calculated energy is 33.49 nJ, which is quite close to the measured result (28.98 nJ).

4) *Stage 4 (Current Ringing,  $t_4 \sim t_5$ ):* The 2DEG channel is fully conducted.  $S_L$  enters the ohmic region and becomes resistive in this stage [30].  $I_g$  continues to charge  $C_{gd}$  and  $V_{gs}$  finally reaches the driving voltage  $V_{gg}$ . This stage is also known as current ringing stage because of the resonance caused by  $C_{eq}$  and  $L_{stray}$ . As mentioned earlier, the loss contributed by ringing can be neglected and only the conduction loss  $E_{cond}$  is considered in this stage.

## B. Loss During Turn-OFF Transition

Similar to the turn-ON transition, the turn-OFF transition can be subdivided into four stages: 1) the turn-OFF delay, 2) the voltage commutation, 3) the current commutation, and 4) the voltage ringing stages. The loss during this transition is normally smaller than that of turn-ON transition, and can be classified as soft-switching [31]. The mechanism of losses  $E_{vi,off}$ ,  $E_{oss}$ , and  $E_{sd}$  during turn-OFF transition will be discussed as follows.

1) *Stage 5 (Turn-OFF Delay,  $t_6 \sim t_7$ ):* When  $V_{ctrl}$  becomes low,  $C_{gs}$  discharges and  $V_{gs}$  begins to fall. At the beginning of this stage,  $S_L$  is still in the ohmic region and its 2DEG channel is fully conducted.  $I_d$  and  $V_{ds}$  can be regarded as nearly constant, and therefore, this stage contributes to  $E_{cond}$  only.

2) *Stage 6 (Voltage Commutation,  $t_7 \sim t_8$ ):*  $V_{gs}$  continues to fall and reaches the Miller plateau.  $S_L$  enters into the saturation region. In contrast to stage 3,  $C_{eq}$  is charging, whereas  $C_{oss1}$  is discharging, as shown in Fig. 3(b).  $V_{ds}$ ,  $I_d$ , and  $I_{ch}$  are expressed, respectively, as follows:

$$V_{ds} = V_{dc} + V_{ds1} + L_{stray} \frac{dI_d}{dt} \quad (7)$$

$$I_d = I_L - C_{oss1} \frac{dV_{ds1}}{dt} \quad (8)$$

$$I_{ch} = I_d - C_{eq} \frac{dV_{ds}}{dt}. \quad (9)$$

It is found that  $C_{\text{oss}1}$  leads to a current knee and, thus, a reduction of turn-OFF  $V - I$  overlapping loss  $E_{\text{vi,off}}$ . Since the charging current of  $C_{\text{eq}}$  is from the external current source instead of the 2DEG channel, the power dissipation is not contributed by the device. This stage ends when  $S_H$  starts to conduct reversely. In summary, this stage contribute to  $E_{\text{vi,off}}$ .

3) *Stage 7 (Current Commutation,  $t_8 \sim t_9$ ):* As  $V_{\text{ds}}$  increases beyond  $V_{\text{dc}} + V_{\text{sd}}$  due to the effect of stray inductance and reverse voltage drop of  $S_H$ ,  $C_{\text{eq}}$  continues to be charged by  $I_d$ . Thus,  $I_d$  continues to fall to zero with the correlation shown in (1), and  $I_{\text{ch}}$  is the same as (9) and  $V_{\text{ds}}$  is given by

$$V_{\text{ds}} = V_{\text{dc}} + V_{\text{sd}} + L_{\text{stray}} \frac{dI_d}{dt}. \quad (10)$$

Losses in this stage contribute to  $E_{\text{sd}}$  and  $E_{\text{vi,off}}$ . At the end of stage 7,  $I_{\text{ch}}$  and  $I_d$  drops to zero since  $V_{\text{gs}} = V_{\text{th}}$ . Meanwhile,  $V_{\text{ds}}$  reaches to its peak. This relationship is useful to locate the alignment point to calibrate the propagation delay of voltage and current probes, and will be discussed in the following section.

Similar to the analysis in stage 3, the current commutation in this stage can be neglected under the extreme condition of low load current and large  $C_{\text{eq}}$  as explained by (8) and (9),  $I_{\text{ch}}$  may drop to nearly zero in stage 6. Thus, there is no  $E_{\text{vi,off}}$  in stage 7 as there is no  $V - I$  overlapping. The analysis agrees well with the previous work in [32], where the turn-OFF transient waveforms are divided into three cases.

4) *Stage 8 (Voltage Ringing,  $t_9 \sim t_{10}$ ):* In this stage,  $S_L$  is in the cut-off region due to the continuous discharge of  $C_{\text{gs}}$ , thus  $V_{\text{gs}}$  is below  $V_{\text{th}}$ . This stage is also known as voltage ringing stage due to the resonance caused by  $C_{\text{eq}}$  and  $L_{\text{stray}}$ . Compared with the current ringing in stage 4, the voltage ringing is more significant since it is less damped with 2DEG channel open, which also leads to significant current ringing with the expression given by

$$I_d = C_{\text{eq}} \frac{dV_{\text{ds}}}{dt}. \quad (11)$$

It can be found the phase of  $V_{\text{ds}}$  lags that of  $I_d$  by  $90^\circ$ .

### C. Switching Loss in Measurement

Based on the earlier analysis, it can be concluded that the theoretical total loss  $E_{\text{tt}}$  is given by

$$E_{\text{tt}} = E_{\text{ton}} + E_{\text{toff}} \quad (12)$$

where the theoretical turn-ON and turn-OFF losses,  $E_{\text{ton}}$  and  $E_{\text{toff}}$ , are due to the overlapping between the waveforms of  $V_{\text{ds}}$  and  $I_{\text{ch}}$ , and are expressed as

$$\begin{cases} E_{\text{ton}} = \int_{t_2}^{t_4} V_{\text{ds}} I_{\text{ch}} dt = E_{\text{vi,on}} + E_{\text{oss}} \\ E_{\text{toff}} = \int_{t_7}^{t_9} V_{\text{ds}} I_{\text{ch}} dt = E_{\text{vi,off}} \end{cases} \quad (13)$$

and  $E_{\text{oss}}$  can be derived from the energy balance theory that the energy stored in the turn-OFF transition will be fully dissipated in the turn-ON transition through the 2DEG channel of GaN HEMT, as given by

$$E_{\text{oss}} = \int_0^{V_{\text{dc}} + V_{\text{sd}}} V_{\text{ds}} C_{\text{oss}} dV_{\text{ds}}. \quad (14)$$

As measuring  $I_{\text{ch}}$  is impossible and only  $I_d$  can be measured in the DPT experiment, which leads to an error in the switching loss. The measured turn-ON and turn-OFF losses,  $E_{\text{mon}}$  and  $E_{\text{moff}}$ , are given by

$$\begin{cases} E_{\text{mon}} = E_{\text{vi,on}} \\ E_{\text{moff}} = E_{\text{vi,off}} + E_{\text{oss}} \end{cases} \quad (15)$$

and the measured total loss  $E_{\text{mt}}$  is given by

$$E_{\text{mt}} = E_{\text{mon}} + E_{\text{moff}} = E_{\text{tt}}. \quad (16)$$

Based on the previous equations, it can be found that  $E_{\text{oss}}$  is measured as part of  $E_{\text{moff}}$ , whereas it is part of  $E_{\text{ton}}$ . It is because that the discharging current of  $C_{\text{oss}}$  and  $C_p$  flows through the 2DEG channel and cannot be measured. Fortunately, the same amount of  $E_{\text{oss}}$  can be measured in the turn-OFF transition. Thus, the measured total loss is the same as that of the theoretical derivation. For simplicity, in the rest of this article,  $E_{\text{mon}}$  and  $E_{\text{moff}}$  are denoted as  $E_{\text{on}}$  and  $E_{\text{off}}$ , respectively.

## III. IMPACT OF PCB PARASITIC ELEMENTS ON SWITCHING LOSS

The fast switching speed makes GaN HEMTs highly sensitive to the PCB parasitic elements, which is the cause of most electromagnetic interference (EMI) issues. The ringing affects the measured waveforms and if the voltage spike at turn-OFF is high enough, it may even lead to the avalanche breakdown. In addition, as the switching loss is estimated based on  $V - I$  overlapping, it may also influence the accuracy of the estimated switching loss. Therefore, this section proposes a modified SPICE model to include the influence of  $L_{\text{stray}}$  and  $C_p$  in the loss estimation. Based on the proposed model, the relationship between the parasitic elements and the switching loss can be analyzed with the parasitic elements as varying parameters.

### A. Modified SPICE Model

To assess the influence of parasitic elements, a SPICE model of GaN HEMT EPC2014 is modified based on the device model from manufacturer [33]. In this model, the relationship between  $V_{\text{ds}}$  and  $I_d$  is given by

$$I_d = \frac{k_1 \ln(1 + \exp((V_{\text{gs}} - k_2)/k_3)) V_{\text{ds}}}{1 + \max((k_4 + V_{\text{gs}}), k_5) V_{\text{ds}}} \quad (17)$$

where parameters  $k_x$  ( $x = 1, 2, \dots, 5$ ) can be obtained through curve fitting with the electrical characteristic of the device given in the datasheet [34]. The three junction capacitances ( $C_{\text{gs}}$ ,  $C_{\text{gd}}$ ,  $C_{\text{ds}}$ ) are voltage-dependent and can be expressed in terms of the following function:

$$C_x(V) = P_{1x} \ln(P_{2x} + \exp((V + P_{3x})/P_{4x})) + P_{5x}. \quad (18)$$

With (18), the three junction capacitances are described by

$$\begin{cases} C_{\text{gs}} = C_{a1}(V_{\text{gs}}) + C_{a2}(V_{\text{ds}}) \\ C_{\text{gd}} = C_{b1}(V_{\text{gs}}) + C_{b2}(V_{\text{gd}}) + C_{b3}(V_{\text{gd}}) \\ C_{\text{ds}} = C_{c1}(V_{\text{ds}}) + C_{c2}(V_{\text{ds}}) \end{cases} \quad (19)$$

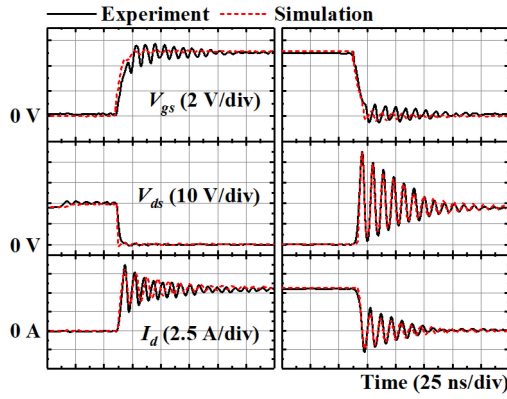


Fig. 4. Switching waveforms comparison between experiment (black solid line) and simulation (red dashed line) when  $L_{loop} = 12.25$  nH,  $C_p = 1$  pF.

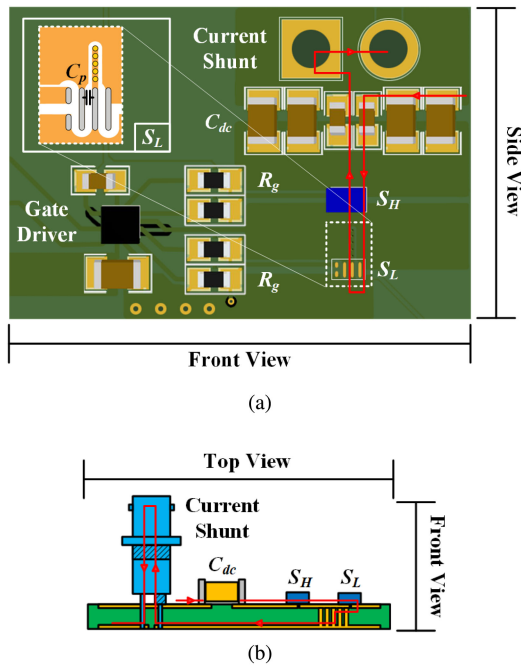


Fig. 5. Low-parasitic PCB design and power loop (red line) of DPT circuit for GaN HEMT. (a) Top view. (b) Side view.

where the respective terms can be obtained from  $C - V$  curve of the datasheet.

To validate the modified SPICE model, it is tested in a DPT circuit with the configuration same as the experiment ( $R_g = 10 \Omega$ ,  $V_{dc} = 18$  V,  $I_L = 6$  A, and  $L = 31 \mu\text{H}$ ).  $L_{loop}$  is extracted from (20) and found to be 12.25 nH and  $C_p$  is extracted by FEM and found to be 1 pF. The comparisons between simulated and measured switching waveforms of  $V_{gs}$ ,  $V_{ds}$ , and  $I_d$  are shown in Fig. 4. It is found that the measured and simulated results show good agreement.

### B. Impact of Stray Inductance

The total stray inductance in a DPT circuit consists of packaging inductances of GaN device, decoupling capacitor, and the parasitic inductance of PCB, as shown in Fig. 5. With the advanced packaging technology such as land grid array [35],

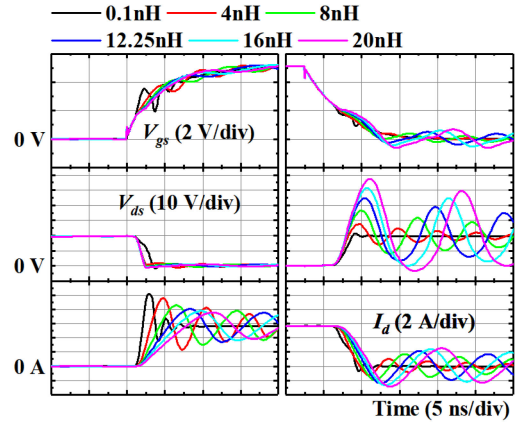


Fig. 6. Simulated switching waveforms with variation of  $L_{loop}$  (0.1–20 nH) at  $C_p = 1$  pF.

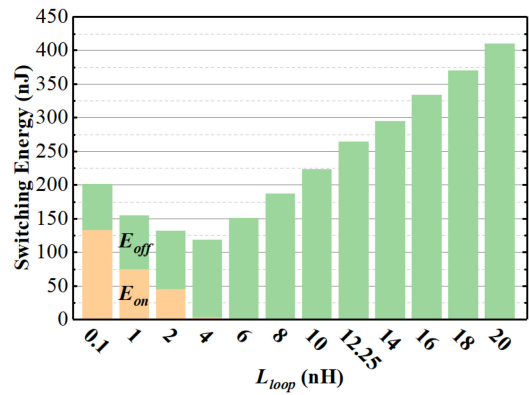


Fig. 7. Simulated switching losses during turn-ON and turn-OFF transitions with varying  $L_{loop}$  (0.1–20 nH) at  $C_p = 1$  pF.

the packaging inductance can be as low as 0.1 nH. The packaging inductance contributed by the decoupling capacitor can be minimized by using several surface mounted device ceramic capacitors in parallel. Thus, the PCB parasitic inductance  $L_{loop}$  dominates  $L_{stray}$ . The value of  $L_{loop}$  can be estimated from the resonant frequency at the switching transition given by

$$L_{loop} = \frac{1}{C_{eq}(2\pi f_s)^2} - L_{pkg} \quad (20)$$

where the total packaging inductance  $L_{pkg}$  is estimated to be 0.2 nH.

The influence of different values of  $L_{loop}$  on the switching loss is further investigated by varying it from 0.1 to 20 nH to cover good to bad PCB layout design. As expected, Fig. 6 shows that higher  $L_{loop}$  has a strong impact on the switching waveforms, which decreases  $di/dt$  and increases  $dv/dt$  at both turn-ON and turn-OFF transitions. During the turn-ON transition,  $V_{ds}$  directly drops to zero before  $I_d$  rise to  $I_L$ , which means the stage 3 can be partially or fully eliminated. Thus,  $E_{on}$  decreases with the increase of  $L_{loop}$ . During the turn-OFF transition, the fast voltage slew rate and slow current slew rate, together with the high-voltage spike, increase  $E_{off}$  significantly.

Fig. 7 shows the breakdown of switching losses during turn-ON and turn-OFF with varying  $L_{loop}$ . When  $L_{loop} = 0.1$  nH,  $E_{on}$  is 133.5 nJ, which is larger than  $E_{off} = 67.3$  nJ. With higher

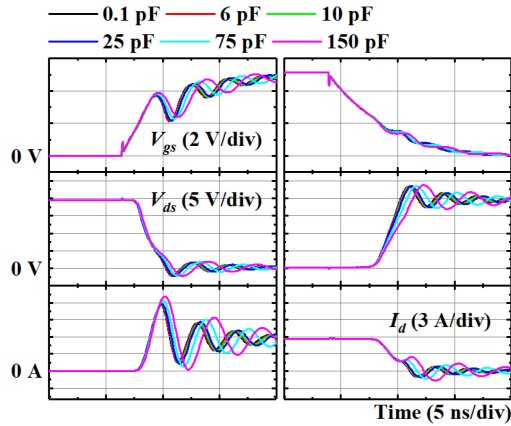


Fig. 8. Simulated switching waveforms with varying  $C_p$  (0.1–150 pF) at  $L_{loop} = 1$  nH.

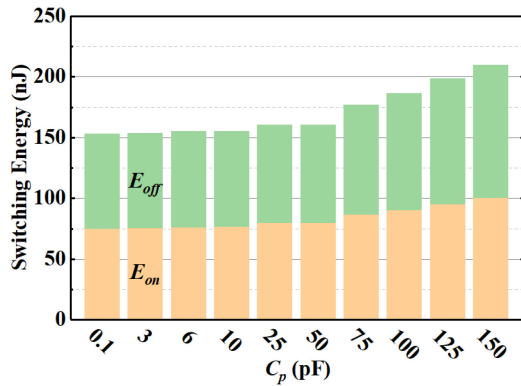


Fig. 9. Simulated switching losses during turn-ON and turn-OFF transitions with varying  $C_p$  (0.1–150 pF) at  $L_{loop} = 1$  nH.

$L_{loop}$ ,  $E_{off}$  further increases and  $E_{on}$  becomes negligible, which agrees well with the findings in [4]. Based on Fig. 7, it reveals that there is an optimal  $L_{loop}$  to achieve minimum total switching loss and highest conversion efficiency. Hence, the lowest  $L_{loop}$  is not always desired, especially for soft-switching application.

### C. Impact of Parasitic Capacitance

With the increase of switching frequency and more PCB layers, loss from the parasitic capacitance can no longer be neglected [36]. As shown in Fig. 1,  $C_p$  is antiparallel with  $S_L$ , thus it requires additional current to complete the charging and discharging processes. This additional loss contributes to  $E_{oss}$  and it has to be evaluated separately.

By varying  $C_p$  from 0.1 to 150 pF, its influence on the switching waveforms is investigated, as shown in Fig. 8. Higher  $C_p$  increases  $di/dt$  and decreases  $dv/dt$ . During the turn-ON transition,  $I_d$  rises to  $I_L$  before  $V_{ds}$  drops to zero, thus  $E_{on}$  increases with the increase of  $C_p$ . During the turn-OFF transition,  $I_d$  drops at first and then rises up. The reason for the changes on  $I_d$  is due to the charging of  $C_{oss}$  and  $C_p$ . This current bump results in an increase of  $E_{off}$ .

The breakdown of switching losses during turn-ON and turn-OFF transitions with varying  $C_p$  is shown in Fig. 9. When  $C_p$  is 0.1 pF,  $E_{on}$  is 75.5 nJ, and  $E_{off}$  is 77.6 nJ. When  $C_p$  is 150

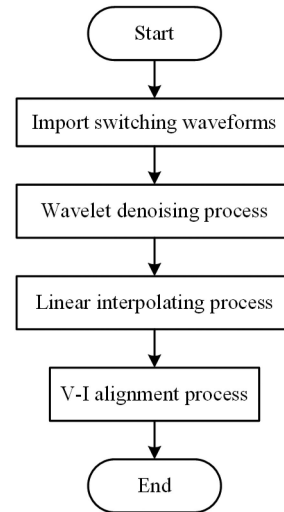


Fig. 10. Flowchart of postprocessing deskew method.

pF, which is close to  $C_{oss}$ ,  $E_{on}$  is 100.8 nJ, and  $E_{off}$  is 108.7 nJ. Further increase in  $C_p$ , the incremental changes of both  $E_{on}$  and  $E_{off}$  are almost the same. Interestingly, it is found that even when  $C_p$  increases to 10 pF, the switching loss shows a slight difference compared with that at 0.1 pF. Hence, the influence of  $C_p$  can be neglected for the 2-layer PCB design in this article.

## IV. POSTPROCESSING-TECHNIQUE-BASED SWITCHING LOSS ESTIMATION METHOD

This section presents a postprocessing method to deskew switching waveforms, which are obtained from the DPT experiment. First, the test setup is introduced and issues that can impact the measurement accuracy are discussed. Then, the mechanism of postprocessing technique is illustrated, with the flowchart shown in Fig. 10. To better understand the proposed method, three processes:  $V - I$  alignment, wavelet denoising, and linear interpolating processes are described in accordance with the sequence shown in Fig. 10. The  $V - I$  alignment process based on the switching characteristics discussed in Section II will be described first. Then, the influence of oscilloscope background noise on the loss estimation is minimized based on a wavelet denoising technique. Finally, linear interpolating algorithm is applied to further reduce the alignment error due to the limitation of sampling rate.

### A. Basic Principle

The DPT method is the most widely adopted measurement setup to estimate the switching loss. Differing from the conventional DPT, which uses a chopper circuit, the synchronous DPT is a half-bridge (HB) configuration with two complementary gate driving signals to simulate the synchronous operating mode, as shown in Fig. 1. The value of main components including the load inductance  $L$  and the dc-link capacitance  $C_{dc}$  must be carefully chosen to ensure that the dc-link voltage  $V_{dc}$  and load current  $I_L$  are constant during the switching transition. The gate driver circuit has to be carefully designed to capitalize the high-speed capability of GaN devices. These designs are crucial

for the DPT circuit and have been well reported and will not be repeated here [20]. In the layout design, the vertical loop (red line in Fig. 5) is adopted to benefit from magnetic flux self-cancellation to reduce  $L_{\text{loop}}$  [37].

Oscilloscope and probe selections are equally important, as the bandwidth of the measurement system ( $\text{BW}_{-3\text{dB}}$ ) is directly affected by the oscilloscope  $\text{BW}_{-3\text{dB}(S_p)}$  and the probe  $\text{BW}_{-3\text{dB}(P_b)}$ , the system bandwidth is governed by

$$\text{BW}_{-3\text{dB}} = \frac{1}{\sqrt{\frac{1}{\text{BW}_{-3\text{dB}(S_p)}^2} + \frac{1}{\text{BW}_{-3\text{dB}(P_b)}^2}}}. \quad (21)$$

Due to the fast switching speed of GaN HEMT, any propagation delay of the measurement system can contribute toward measurement errors. Usually, it can be minimized with a short cable [38].

Besides, the parasitic elements introduced by the probes have to be taken into account for the measurement of GaN HEMT. The voltage probe with the lowest possible coupling capacitance and shortest ground lead length should be selected. The current measurement is usually the challenge when comparing with the voltage measurement, as the current probe has to be inserted into the power loop, which increases  $L_{\text{loop}}$ . The conventional current transformer and Rogowski-coil are not suitable due to their relatively large size. The current surface probe method [39] and integrated current measurement method [40] have superior performance but at the expense of complexity and higher costs. As a compromise, the current shunt is chosen for the experiment.

### B. $V - I$ Alignment Process

Considering the high switching speed of GaN HEMTs, any slight propagation delay between the voltage and current probes can contribute to large measurement error. Although the oscilloscope provides the sampling synchronization between voltage probes (refer to Fig. 13), the synchronization between voltage and current probes are challenging without the use of external deskewing. Thus, the need for postprocessing alignment is required.

The first step of this process is to locate a suitable point based on the switching mechanism. Based on the discussion in the Section II,  $t_9$  (i.e., end of stage 7 and beginning of stage 8) is selected as the alignment point, as shown in Fig. 2. At this point,  $V_{\text{ds}}$  reaches the peak value, thus  $dV_{\text{ds}}/dt = 0$ . From (9), it can be found that  $I_{\text{ch}} = I_d$ . Considering the phase of  $V_{\text{ds}}$  that lags that of  $I_d$  by  $90^\circ$  during the voltage ringing stage, it can be referred that  $I_{\text{ch}} = I_d = 0$  at  $t_9$ . Thus,  $V_{\text{gs}}$  should be equal to  $V_{\text{th}}$  so that the 2DEG channel is in cut-off. Bearing in mind that  $V_{\text{th}}$  is temperature-dependent, the alignment point for  $V_{\text{gs}}$  should take the temperature effect into consideration, as shown in Fig. 11. To summarize, the alignment point is selected to be  $t_9$  with  $V_{\text{ds}} = V_{\text{dsmax}}$ ,  $V_{\text{gs}} = V_{\text{th}}$ , and  $I_d = 0$ . The alignment point  $t_9$  is complementary to  $t_2$  at the turn-ON transition. However,  $t_9$  is much easier to be located since the voltage and current ringings are inevitable due to the effect of  $L_{\text{stray}}$ .

The  $V - I$  alignment process can be implemented in MATLAB with its flowchart shown in Fig. 12. First, the three measured signal data  $x(n) = [V_{\text{ds}}, V_{\text{gs}}, I_d]^T$  are imported. Then, the

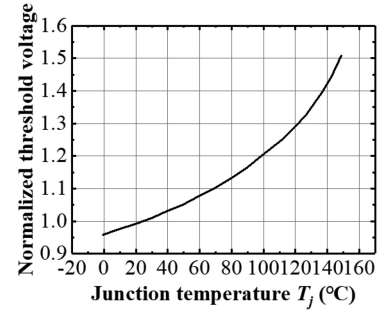


Fig. 11. Relationship between  $V_{\text{th}}$  and junction temperature  $T_j$  of GaN HEMT EPC2014 when  $I_d$  is 2 mA.

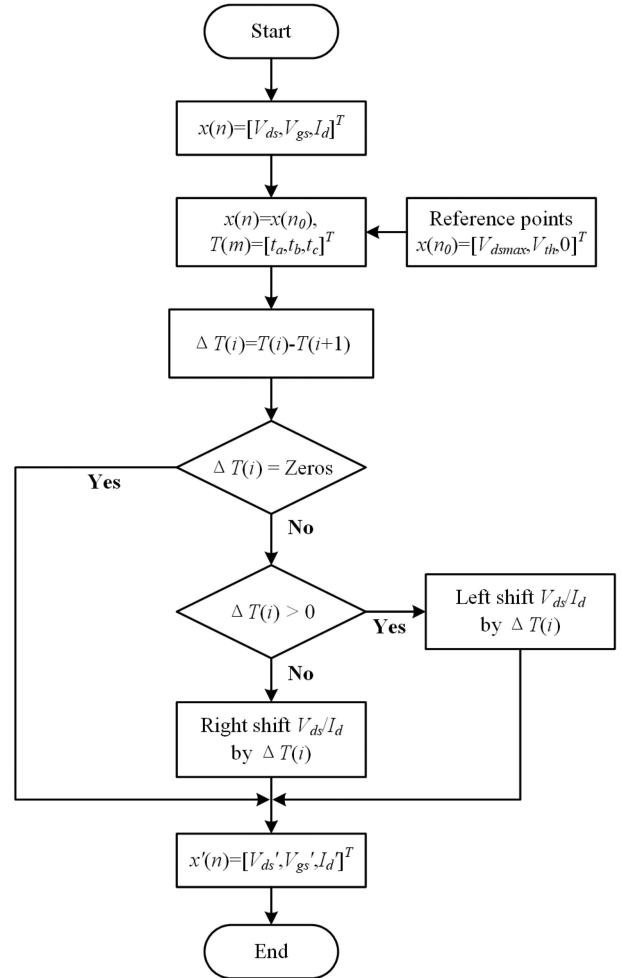


Fig. 12. Flowchart of  $V - I$  alignment process.

reference points  $t_a$ ,  $t_b$ , and  $t_c$  at  $V_{\text{ds}} = V_{\text{dsmax}}$ ,  $V_{\text{gs}} = V_{\text{th}}$ , and  $I_d = 0$  are obtained, respectively. By comparing these reference points, the time differences can be determined as

$$\Delta T(i) = T(i) - T(i+1) \quad (22)$$

where  $i$  is chosen as 1 or 2 for representation of  $V_{\text{ds}} - V_{\text{gs}}$  or  $V_{\text{gs}} - I_d$  probe alignment. If  $\Delta T(i) > 0$ , it indicates that  $V_{\text{ds}}$  ( $I_d$ ) probe has a propagation delay and it should be shifted left with respect to  $V_{\text{gs}}$  by the time difference  $\Delta T(i)$ . Similarly,

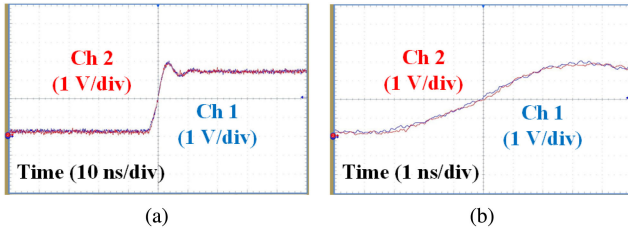


Fig. 13. Background noise on measured waveforms observed on the oscilloscope at different time scale using the voltage probe TDP1000 at channel 1 and 2. (a) 10 ns/div. (b) 1 ns/div.

it should be shifted right if  $\Delta T(i) < 0$ . The process continues till  $\Delta T(i) = 0$  and finally, the alignment of the signals  $x(n) = [V'_{ds}, V'_{gs}, I'_d]^T$  is obtained.

### C. Wavelet Denoising Process

The measured signals from the digital oscilloscope are always mixed with noises. Fig. 13 illustrates the undesirable random zig-zag noises mixed with the desired voltage signals. The noise is contributed by the front-end analog conditioning circuit thermal noise, quantization noise in analog-to-digital converter (ADC), etc. Since the proposed loss estimation method is based on the measured waveforms, the noise may not only make the  $V - I$  alignment more difficult but also affect the accuracy of the loss estimation. Therefore, the influence of the noise must be reduced. These noises can be modeled as Gaussian white noises, which can be reduced by the wavelet denoising technique. For the following discussion, the voltage and current signal, which are the wanted signals, will be denoted as “signal,” whereas the Gaussian white noise, which is unwanted noise, will be denoted as “noise.”

The denoising process is based on the wavelet transform, which overcomes shortcomings of Fourier transform. The wavelet transform uses a set of basis analysis functions known as “wavelet,” and then provide a mother wavelet  $\phi(t)$ . By shifting and scaling, one can get a set of basis functions called as “wavelet basis,” as given by

$$\phi_{a,b}(t) = \frac{1}{\sqrt{a}} \phi\left(\frac{t-b}{a}\right). \quad (23)$$

The wavelet transform is the inner product of input signal  $x(t)$  and  $\phi_{a,b}(t)$

$$\text{WT}_x(a, b) = \int x(t) \phi_{a,b}^*(t) dt = \langle x(t), \phi_{a,b}(t) \rangle \quad (24)$$

where  $a$  is positive and defines the scale, whereas  $b$  is any real number and defines the shift. Therefore, based on the wavelet transform, the denoising process can be applied with the flowchart shown in Fig. 14. The measured waveforms  $x(n) = [V_{ds}, V_{gs}, I_d]^T$  are defined as original waveforms

$$x(n) = s(n) + u(n) \quad (25)$$

where  $s(n)$  is the signal, and  $u(n)$  is the noise. By applying the wavelet transform on both sides of (25), it results in

$$\text{WT}_x(a, b) = \text{WT}_s(a, b) + \text{WT}_u(a, b). \quad (26)$$

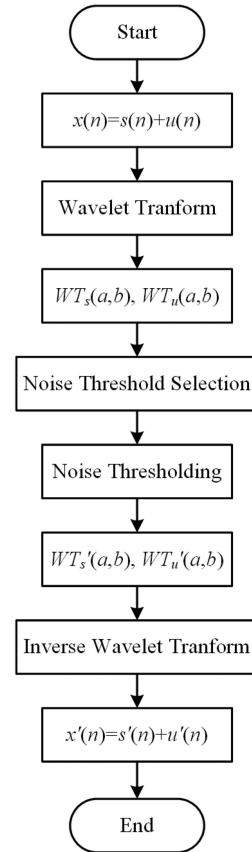


Fig. 14. Flowchart of wavelet denoising method.

Since  $u(n)$  is Gaussian white noise, it still remains as the Gaussian white noise after the wavelet transform based on the signal processing theory. Therefore, the correlation of  $s(n)$  is removed, and its energy is only on the wavelet coefficient, which is the maximum value of modules in all scales. Noting that  $u(n)$  is still the Gaussian white noise, its wavelet coefficient is uncorrelated and distributed in each time axis of all scales. Hence, the wavelet denoising process is to keep the large coefficient unaltered and set the small coefficient to zero. Then, by applying the inverse wavelet transform, the noise can be minimized.

In reality, the small coefficient cannot be simply set to zero, as there is useful information in it. Therefore, the wavelet shrinkage is applied and the soft thresholding is selected as

$$\omega_\lambda = \begin{cases} [\text{sign}(\omega)] (|\omega| - \omega), & |\omega| \geq \omega \\ 0, & |\omega| < \omega \end{cases} \quad (27)$$

where  $\omega_\lambda$  is the wavelet coefficients value imposed threshold,  $\lambda$  is the threshold value, and  $\omega$  is the wavelet coefficient. Therefore, the key issue is how to adjust threshold to obtain the optimal threshold for the signal. To overcome this issue, the principle of Stein's unbiased risk estimate is deployed as follows:

$$\begin{cases} sx2(k) = (\text{SRT}_x)^2 \\ \lambda = \sqrt{sx2(k_{\min})} \end{cases}, \quad k = 0, 1, \dots, N-1 \quad (28)$$

where  $N$  is the length of signal array,  $sx2(k)$  is a new signal array,  $\text{SRT}_x$  is the array of absolute value of  $x(n)$  sorted from

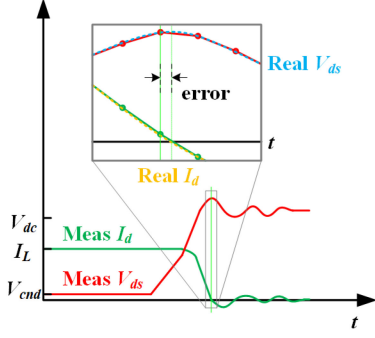


Fig. 15. Switching waveforms of  $V_{ds}$  and  $I_d$  in turn-OFF transition show how the fixed sampling step of oscilloscope can bring in the error.

ascending order, and  $k_{\min}$  is the optimal point, which is given by

$$\min \left[ 1 - \frac{2k}{N} + \frac{1}{N} \sum_{j=1}^k sx2(j) + \left( 1 - \frac{k}{N} \right) sx2(N - k) \right]. \quad (29)$$

Both (28) and (29) are implemented in the MATLAB wavelet Toolbox so that the noise can be effectively filtered and minimized.

#### D. Linear Interpolation Process

Even the sampling rate of advanced digital oscilloscope is in the range of 1–100 GS/s, it is still a challenging task to align precisely at the point when  $V_{ds} = V_{ds\max}$ ,  $V_{gs} = V_{th}$ , and  $I_d = 0$ . As shown in Fig. 15, the GaN HEMT EPC2014 has a threshold voltage  $V_{th} = 1.4$  V. However, the nearest point to  $V_{th}$  is 1.43 V, which can lead to measurement error due to inability of precise alignment.

The measured signal  $x(n)$  is sampled at a fixed sampling step of  $t_s$ . As  $t_s$  is usually small enough (0.2 ns for 5 GS/s in this work), the interpolation point between the two sampling points can be estimated by the linear approximation. Hence, the linear interpolation method is adopted and the linear interpolated data  $x'(n)$  can be given by

$$x'(n) = x(k) + \frac{(x(k+1) - x(k))(t'(n) - t(k))}{t(k+1) - t(k)} \quad (30)$$

where  $t(k+1) - t(k)$  is constant as  $t_s$  is a fixed value. It should be noticed that this process is before the  $V - I$  alignment process, and therefore, its influence on the alignment has to be considered. For example, the time difference between  $V_{gs} = V_{th}$  and  $I_d = 0$  for one group of measured waveforms is 0.06 ns, whereas that for  $V_{ds} = V_{ds\max}$  and  $V_{gs} = V_{th}$  is 0.13 ns. The most straightforward way is to interpolate data for each waveform using the time steps of 0.06 and 0.13 ns. However, by doing so, it makes the sampling rate of the measured waveforms a variable, which can lead to distorted waveforms during the  $V - I$  alignment process. Thus, it is better to interpolate all data by a fixed time step of  $t_{ip}$ , which must be carefully chosen.

It is a balance between smaller time step and longer computation time. In order to select an appropriate  $t_{ip}$ , three time points  $t_a$ ,  $t_b$ , and  $t_c$  are defined, which are the nearest sampling point to  $V_{ds} = V_{ds\max}$ ,  $V_{gs} = V_{th}$ , and  $I_d = 0$ , respectively. To

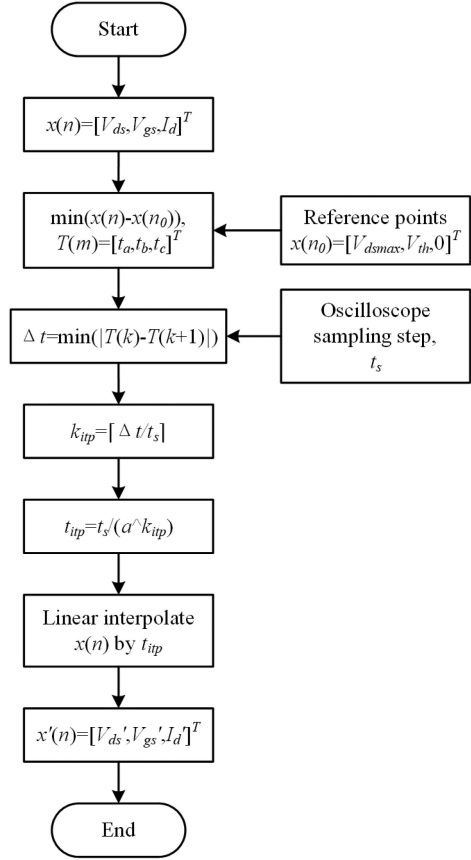


Fig. 16. Flowchart of linear interpolation process.

simplify the analysis,  $t_b$  is taken as the reference point, and the time difference  $\Delta t$  is given by

$$\Delta t = \min(|T(k) - T(k+1)|) \quad (31)$$

where  $T(m) = [t_a, t_b, t_c]^T$ ,  $k$  is chosen as either 1 or 2 to represent the time difference of  $V_{ds} - V_{gs}$  or  $V_{gs} - I_d$ , respectively. Then compared with the original  $t_s$ , the interpolation rate  $k_{ip}$  for each data interval is obtained as follows:

$$k_{ip} = \frac{\Delta t}{\Delta t_s}. \quad (32)$$

Finally, the optimized  $t_{ip}$  is obtained by

$$t_{ip} = \frac{t_s}{a^{\lceil k_{ip} \rceil}} \quad (33)$$

where the value of  $k_{ip}$  is rounded up to an integer to avoid a rational number, and  $a$  is an integer number, which is chosen as 10 in this article. The flowchart of linear interpolation process is given in Fig. 16.

#### V. EXPERIMENTAL SETUP

Using the GaN HEMT EPC2014 (40 V/10 A) as the device under test, DPT experimental data are collected to validate the proposed loss estimation method. The HB GaN driver LMG1205 from Texas Instruments is chosen to drive GaN HEMT. The external gate resistances for both turn-ON and turn-OFF transitions are 10  $\Omega$ . The dead time  $T_d$  is set as 50 ns to avoid the arm

TABLE I  
SPECIFICATIONS OF THE DPT EXPERIMENT

Parameter	Symbol	Value
DC bus voltage	$V_{dc}$	18 V
Decoupling capacitance	$C_{dc}$	40.2 $\mu$ F
Load inductance	$L$	31 $\mu$ H
External gate resistance	$R_g$	10 $\Omega$
Dead time	$T_d$	50 ns

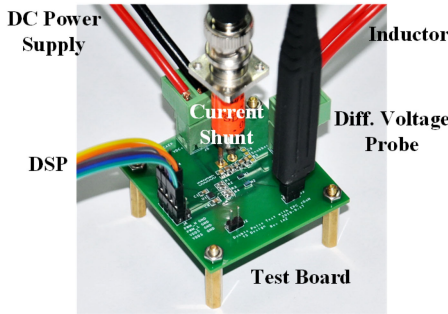


Fig. 17. Experimental setup of DPT for GaN HEMT.

shoot-through. A single-layer, air-core inductor with the value of 31  $\mu$ H is used as the load to minimize the parasitic capacitance as well as to avoid the magnetic saturation. To evaluate the power dissipation at different loadings, the dc voltage is kept constant at 18 V and the load current varies from 2 to 8 A. The specifications of DPT experiment are summarized in Table I, and the experimental setup is shown in Fig. 17.

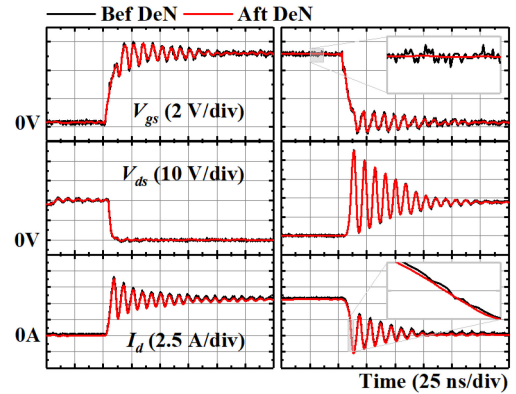
Two wideband differential voltage probes TDP1000 from Tektronix (42 V, 1 GHz) are used to measure the switching waveforms of  $V_{gs}$  and  $V_{ds}$ . For the current measurement, the current shunt SDN-414-025 (bandwidth 1.2 GHz) from T&M Research is used to acquire  $I_d$ . It has a very small rise time of 0.3 ns, which has minimal influence on the current waveform. The experiment is conducted in room temperature (25  $^{\circ}$ C), where  $V_{th} = 1.4$  V according to Fig. 11.

## VI. RESULTS AND DISCUSSION

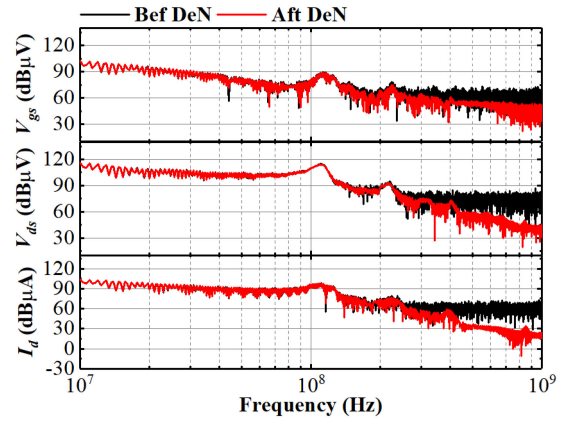
### A. Wavelet Denoising

Since the original measured waveforms from digital oscilloscope contain both signal and noise, the wavelet denoising is applied on the waveforms to filter the background noise. A comparison between the original and denoised waveforms is shown in Fig. 18. As observed in the time domain, the noise is being filtered, whereas the signal remains unaffected. From the frequency-domain results, the low-frequency (<200 MHz) components, which contain the fundamental frequency, high-order harmonics and ringing, are not influenced by the wavelet denoising. In contrast, the high-frequency (>200 MHz) components are being suppressed.

The energies associated with the switching losses, before and after the wavelet denoising process under different load current, are shown in Fig. 19. It is observed that  $E_{on}$  has a random relationship with the current, whereas  $E_{off}$  increases with the current. It can be explained as follows: The switching losses

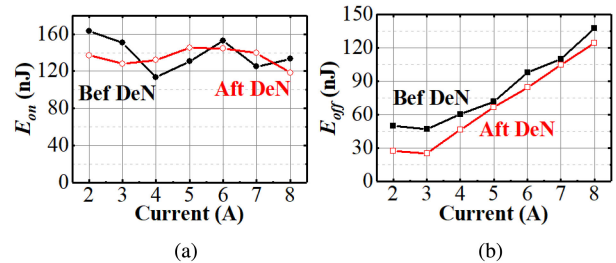


(a)



(b)

Fig. 18. Comparison of switching waveforms before (black line) and after (red line) denoising in (a) time domain and (b) frequency domain.



(a)

(b)

Fig. 19. Comparison of (a)  $E_{on}$  and (b)  $E_{off}$  versus load current before (black line) and after (red line) wavelet denoising process.

have two components, the actual value and the error; and the actual value of  $E_{on}$  is much smaller than the error but that of  $E_{off}$  is the reversed. In addition,  $E_{on} > E_{off}$ , which is because the  $V - I$  alignment has not been conducted. The impact of Gaussian white noise on losses does not show the same tendency due to its random distribution nature.

### B. $V - I$ Alignment

After the wavelet denoising process, the noise component in the signal is being cleaned up. Now the condition with a load current of 6 A is chosen for comparison to observe the effect of the alignment process, as shown Fig. 20. Since  $V_{gs}$  is taken as a reference signal, there are no changes before and after alignment.

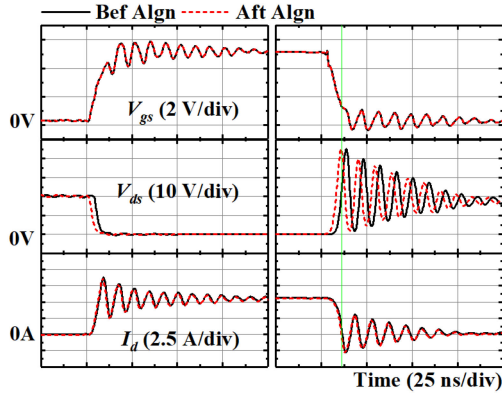


Fig. 20. Comparison of switching waveforms before (black line) and after (red dotted line) alignment with respect to the reference (light green line).

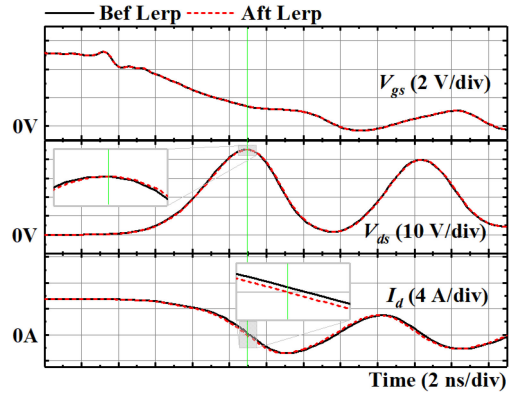


Fig. 22. Comparison of turn-OFF waveforms before (black line) and after (red dotted line) linear interpolation with respect to the reference (light green line).

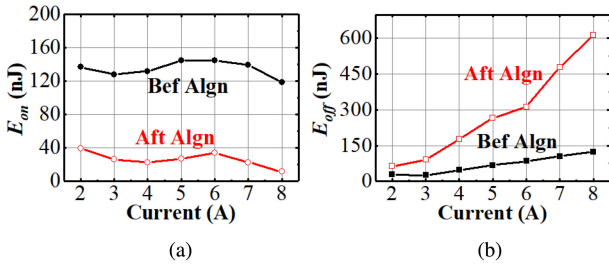


Fig. 21. Comparison of (a)  $E_{on}$  and (b)  $E_{off}$  versus load current before (black line) and after (red line) alignment process.

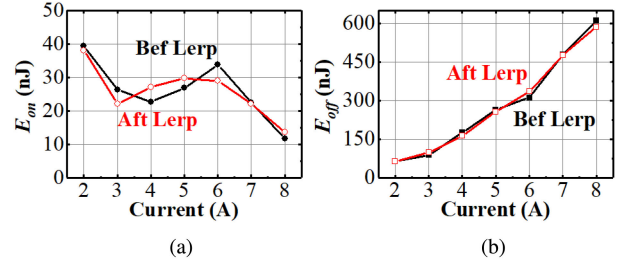


Fig. 23. Comparison of (a)  $E_{on}$  and (b)  $E_{off}$  versus load current before (black line) and after (red line) linear interpolation process.

The light green line in Fig. 20 is to show that the three signals have been correctly aligned. It is clearly observed that  $V_{ds}$  has been moved left, whereas  $I_d$  is moved slightly.

The energies associated with the switching losses before and after the alignment process under different load currents is shown in Fig. 21. It shows that  $E_{on}$  becomes lower, whereas  $E_{off}$  becomes higher, and  $E_{on} < E_{off}$ . It differs from the common knowledge that  $E_{on} > E_{off}$ . For example, when the load current is 6 A,  $E_{on}$  is decreased to 33.7 nJ from 144.6 nJ, whereas  $E_{off}$  is increased to 312.2 nJ from 84.8 nJ. It is due to the use of current shunt, which introduces a stray inductance that affects both  $E_{on}$  and  $E_{off}$ , as discussed in Section III. To summarize, the  $V - I$  alignment process plays a key role of the proposed method. As it is based on the switching mechanism of GaN HEMT, this method can be easily applied.

### C. Linear Interpolation

To further improve the accuracy of alignment process, a linear interpolation process is applied. Fig. 22 shows the turn-OFF waveforms before and after the linear interpolation process. Similar to the previous discussion,  $V_{gs}$  is taken as a reference signal, and a light green line is used as the reference. It is found that  $V_{ds}$  has been shifted to the right slightly, whereas  $I_d$  has been shifted to the left.

The switching energies associated with the switching losses before and after the interpolation process under different load currents are shown in Fig. 23. Comparing with the alignment

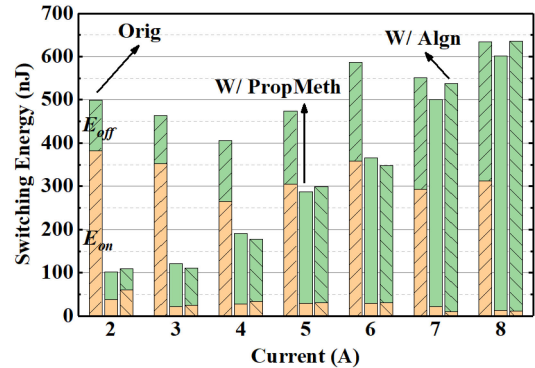


Fig. 24. Switching losses breakdown versus load current without (left) and with (middle) the proposed switching loss estimation method and only with alignment method (right).

process, the impact on the losses is relatively small. For example, when the load current is 6 A,  $E_{off}$  rises to 336.5 nJ from 312.2 nJ, whereas  $E_{on}$  drops to 29.0 nJ from 33.7 nJ. Since the sampling rate of oscilloscope used in this article is 5 GS/s, it has sufficient sampling rate to measure the switching waveforms of GaN HEMT.

### D. Verification of Proposed Switching Loss Estimation Method

Finally, the proposed switching loss estimation method with all the three mentioned processes is implemented. Fig. 24 shows the relationship between the energies associated with the

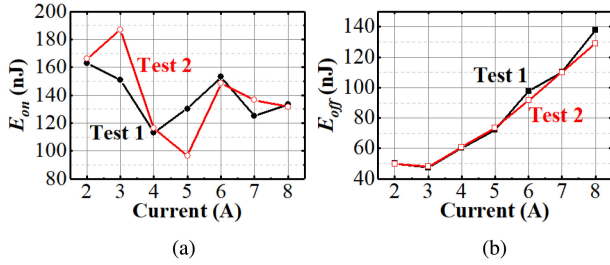


Fig. 25. Original switching energies versus load current for test 1 and test 2. (a)  $E_{on}$ . (b)  $E_{off}$ .

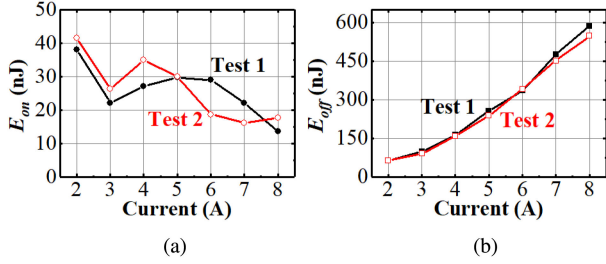


Fig. 26. Calibrated switching energies versus load current for test 1 and test 2 using the proposed method. (a)  $E_{on}$ . (b)  $E_{off}$ .

switching losses versus the load current with and without the proposed method applied. The result with only the alignment process implemented is also presented for better appreciation of the proposed method. It has shown that the original data without the postprocessing method can lead to incorrect information of the losses. For example, the zero-voltage switching (ZVS) is always preferred compared with the zero-current switching (ZCS) in the resonant converter design since  $E_{on} > E_{off}$ . However, for the situation with relatively large stray inductance, the selection of ZVS operation may lead to a lower conversion efficiency. Compared with the wavelet denoising and linear interpolation processes, the  $V - I$  alignment process shows the highest influence on the total loss. Nevertheless, the combination of the three processes has shown a great improvement on the accuracy of the loss estimation.

To check the impact of test environment on the measured results, the proposed method is further applied to two different groups of measurement data, namely test 1 and test 2, respectively. Figs. 25 and 26 show the original and calibrated switching energies versus load current, respectively. It has shown that  $E_{on}$  is slightly different, whereas  $E_{off}$  is nearly the same between the two sets of results. It is expected because the  $V - I$  overlapping for the turn-ON transition is much smaller than that of turn-OFF transition, and the measurement error will have a more significant influence.

## VII. CONCLUSION

Starting with a comprehensive theoretical analysis of the switching mechanism and a careful evaluation of possible impact of PCB parasitic elements on the DPT measurement results, a

postprocessing method has been proposed and described with the objective to estimate the switching loss of GaN HEMT accurately. The proposed method consists of three processes. The first process deals with the wavelet denoising to filter the background noise of the oscilloscope. Then, a  $V - I$  alignment process is developed to reduce errors from the probes' propagation delays. Finally, the linear interpolation process is deployed to further reduce the error from the alignment process due to the possible limitation of the sampling rate. Among the three processes, the  $V - I$  alignment has been demonstrated to be the key factor that has the most significant impact on the accuracy of the estimated switching losses. The proposed switching loss estimation method is experimentally validated on a 40-V/10-A GaN HEMT. The proposed method provides a valuable guideline to analyze and estimate the switching losses of GaN HEMT with confidence in hard- or soft-switching GaN-based converters.

## APPENDIX

Under the low-voltage and large stray inductance condition, only stage 2 is existed, thus  $E_{ton}$  is

$$E_{ton} = \int_{t_2}^{t_3} V_{ds} I_{ch} dt. \quad (A.1)$$

The system equations are

$$\begin{cases} I_d = g_m (V_{gs} - V_{th}) \\ V_{ds} = V_{dc} + V_{sd} - L_{stray} \frac{dI_d}{dt} \\ V_{gs} = V_{gg} - I_g R_g - L_s \frac{dI_d}{dt} \\ I_g = -C_{gd} \frac{dV_{ds}}{dt} \\ I_{ch} = I_d + C_{eq} \frac{dV_{ds}}{dt}. \end{cases} \quad (A.2)$$

Therefore,  $V_{ds}$  and  $I_{ch}$  can be obtained by

$$\begin{cases} V_{ds} = V_{dc} + V_{sd} - L_{stray} g_m \frac{dV_{gs}}{dt} \\ I_{ch} = g_m (V_{gs} - V_{th}) - C_{eq} L_{stray} g_m \frac{d^2 V_{gs}}{dt^2} \\ V_{gs} = V_{gg} - C_{gd} L_{stray} g_m R_g \frac{d^2 V_{gs}}{dt^2} - L_s g_m \frac{dV_{gs}}{dt} \end{cases} \quad (A.3)$$

where  $V_{gs}$ ,  $\frac{dV_{gs}}{dt}$  and  $\frac{d^2 V_{gs}}{dt^2}$  are

$$\begin{cases} V_{gs} = V_{gg} + C_1 e^{-A_1 t} + C_2 e^{-A_2 t} \\ \frac{dV_{gs}}{dt} = -C_1 A_1 e^{-A_1 t} - C_2 A_2 e^{-A_2 t} \\ \frac{d^2 V_{gs}}{dt^2} = C_1 A_1^2 e^{-A_1 t} + C_2 A_2^2 e^{-A_2 t} \end{cases} \quad (A.4)$$

where

$$\begin{cases} A_1 = -\frac{L_s g_m + \sqrt{g_m (g_m L_s^2 - 4 C_{gd} L_{stray} R_g)}}{2 C_{gd} L_{stray} R_g g_m} \\ A_2 = -\frac{L_s g_m - \sqrt{g_m (g_m L_s^2 - 4 C_{gd} L_{stray} R_g)}}{2 C_{gd} L_{stray} R_g g_m} \end{cases} \quad (A.5)$$

$C_1$  and  $C_2$  are constant value, which are related to the initial condition. Therefore,  $E_{\text{ton}}$  is

$$\begin{aligned}
E_{\text{ton}} = & L_{\text{stray}} V_{\text{gg}}^2 g_m^2 (t_2 - t_3) - V_{\text{dc}} V_{\text{gg}} g_m (t_2 - t_3) \\
& + V_{\text{dc}} V_{\text{th}} g_m (t_2 - t_3) - V_{\text{gg}} V_{\text{sd}} g_m (t_2 - t_3) \\
& + V_{\text{sd}} V_{\text{th}} g_m (t_2 - t_3) - L_{\text{stray}} g_m^2 \left( \frac{C_1^2 B_5}{2A_1} + \frac{C_2^2 B_4}{2A_2} \right) \\
& + L_{\text{stray}} V_{\text{gg}} V_{\text{th}} g_m^2 (t_2 - t_3) + \frac{C_1 (V_{\text{dc}} + V_{\text{sd}}) g_m B_2}{A_1} \\
& + \frac{C_2 (V_{\text{dc}} + V_{\text{sd}}) g_m B_1}{A_2} - \frac{2C_1 L_{\text{stray}} (V_{\text{gg}} - V_{\text{th}}) g_m^2 B_2}{A_1} \\
& - \frac{2C_2 L_{\text{stray}} (V_{\text{gg}} - V_{\text{th}}) g_m^2 B_1}{A_2} \\
& + (A_1 C_1^2 B_5 + A_2 C_2^2 B_4) \frac{C_{\text{eq}} L_{\text{stray}}^2 g_m^2}{2} \\
& + ((A_1^2 + A_2^2) L_{\text{stray}} C_{\text{eq}} - 2) \frac{C_1 C_2 L_{\text{stray}} g_m^2 B_3}{A_1 + A_2} \\
& + (L_{\text{stray}} V_{\text{gg}} g_m - V_{\text{dc}}) A_1 C_1 C_{\text{eq}} L_{\text{stray}} g_m B_2 \\
& + (L_{\text{stray}} V_{\text{gg}} g_m - V_{\text{dc}}) A_2 C_2 C_{\text{eq}} L_{\text{stray}} g_m B_1 \\
& - (A_1 C_1 B_2 + A_2 C_2 B_1) C_{\text{eq}} L_{\text{stray}} V_{\text{sd}} g_m
\end{aligned} \tag{A.6}$$

where

$$\begin{cases}
B_1 = e^{-A_2 t_2} - e^{-A_2 t_3} \\
B_2 = e^{-A_1 t_2} - e^{-A_1 t_3} \\
B_3 = e^{-A_1 t_2 - A_2 t_2} - e^{-A_1 t_3 - A_2 t_3} \\
B_4 = e^{-2A_2 t_2} - e^{-2A_2 t_3} \\
B_5 = e^{-2A_1 t_2} - e^{-2A_1 t_3}
\end{cases} \tag{A.7}$$

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